

Measurement of the Superparticle Mass Spectrum in the Long-Lived Stau Scenario at the LHC

Takumi Ito
Tohoku University

In collaboration with R. Kitano and T. Moroi,
arXiv:0910.5853[hep-ph]

特定領域「フレーバー物理の新展開」研究会 2010
2010.2.23

Introduction

Supersymmetry : an important target of the LHC experiments.

Signature of SUSY at the LHC strongly depends on what the LSP in the MSSM sector (= "MSSM-LSP") is!

Popular candidates of MSSM-LSP are:

* The lightest neutralino $\tilde{\chi}_1^0$ ——— Large missing p_T

* **The lightest stau** $\tilde{\tau}_1$

* ...

It may have a long lifetime enough to escape from the detector.
→ we can observe its track.

Introduction

Long-lived stau scenario

↔ Stau behaves as a “heavy muon” in high energy experiments.

* We can measure stau's charge, momentum, velocity and energy.

→ **The stau mass can be measured.** [Ambrosanio et al.(01),
Ellis et al.(06,07)]

→ **All** final state particles (= **SM particles** & stau) in the event
are **visible**. (Of course, except ν .)

* A **slow-moving charged track** (= stau) informs us a **production of superparticles**.

→ we can reduce SM backgrounds significantly.

→ **We have a opportunity to *probe the SUSY in detail!***

Introduction

Long-lived stau scenario

- * It is naturally realized if $\tan\beta$ is large.
- * Cosmological problems can be **avoided** if there is **weekly interacting LSP** (i.e. gravitino, axino, ...).

This scenario is *not an unrealistic scenario.*

→ We may observe *stau's track at the LHC.*

We discuss a determination of the superparticle masses at the LHC experiments in the model with the long-lived stau.

Assumptions and Outline

Assumptions

- * Stau is **stable** in the detector.
- * Stau with $v < 0.9c$ can be identified as “**stau**”, *not muon*. [Ambrosanio et.al.(01)]
- * **No SM backgrounds**. ← Since there is **at least one stau** in the **SUSY event**.
- * The **stau mass has been measured**. (accuracy ~ 100 MeV) [Ambrosanio et al.(01), Ellis et al.(06,07)]

Mass Measurements

1. Neutralino

2. Sleptons

3. Squarks

Charge information of the stau is useful.

The knowledge of the **neutralino mass** helps our reconstruction.

Sample Model

GMSB [Dine,Nelson,Shirman(95)]

$$\Lambda = 60\text{TeV}, M_{\text{mess}} = 900\text{TeV},$$

$$N_5 = 3, \tan \beta = 35, \dots$$

\tilde{g}	1309
\tilde{q}_L	1230
\tilde{q}_R	1180
\tilde{W}^0	426
\tilde{B}	240
\tilde{e}_R	194
$\tilde{\mu}_R$	193
$\tilde{\tau}_1$	149

*Reduced
1 GeV
by hands*

MSSM-LSP

*The decay
products can
be observed.*

masses are in GeV

Monte Carlo Analysis

* Event Gen. by HERWIG6.510

* Fast Detector Sim. by PGS4

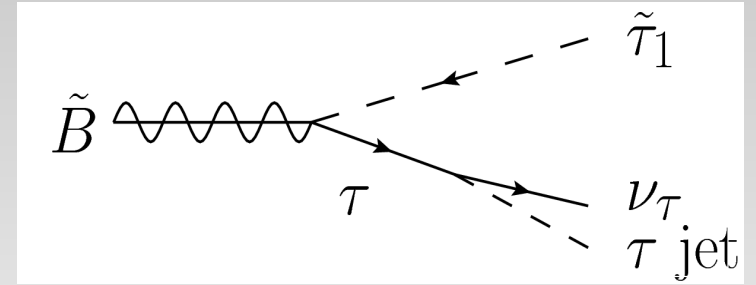
* 67k events are generated.

($\leftrightarrow 100 \text{ fb}^{-1}, 14 \text{ TeV}$)

Neutralino Masses

The lightest neutralino (bino) decay

bino $\rightarrow \tau + \text{stau}$



We consider **hadronic decay** mode of **τ -lepton**.

\leftrightarrow Some part of the energy of τ is taken away by **τ -neutrino**.

$$M_{\tilde{B}} = \sqrt{(p_{\tilde{\tau}_1} + p_{\tau\text{-jet}} + p_{\nu_\tau})^2}$$
$$\geq \sqrt{(p_{\tilde{\tau}_1} + p_{\tau\text{-jet}})^2}$$

Endpoint
 \updownarrow
Bino mass

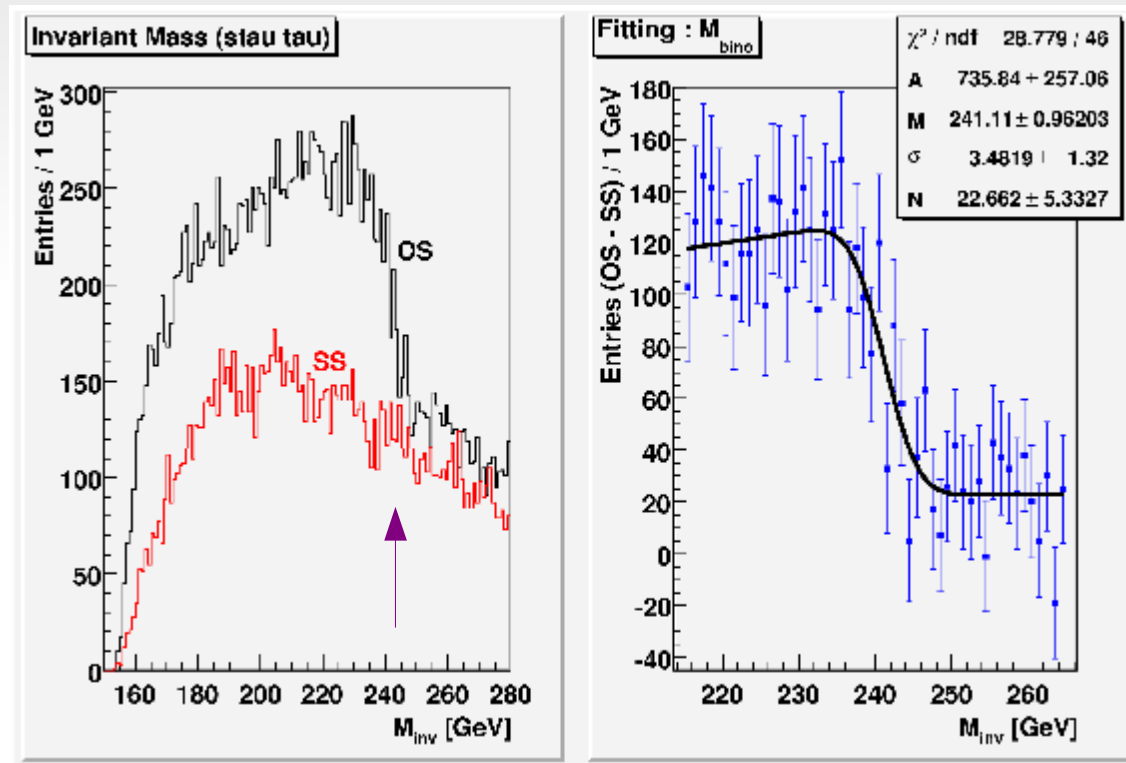
Event Selection

- (1) At least one stau with $0.4c < v < 0.9c$
- (2) At least one τ -tagged jet with $p_T > 15 \text{ GeV}$

Neutralino Masses

To reduce BGs, we adopt a **charge subtraction**:

$$\text{OS } (\tau \text{ stau}) - \text{SS } (\tau \text{ stau}) = \text{Signal}$$



$$M_{\tilde{B}} = 241.1 \pm 1.0 \text{ GeV} \\ (239.5 : \text{underlying})$$

Uncertainties

* Statistics

$$\pm 1 \text{ GeV}$$

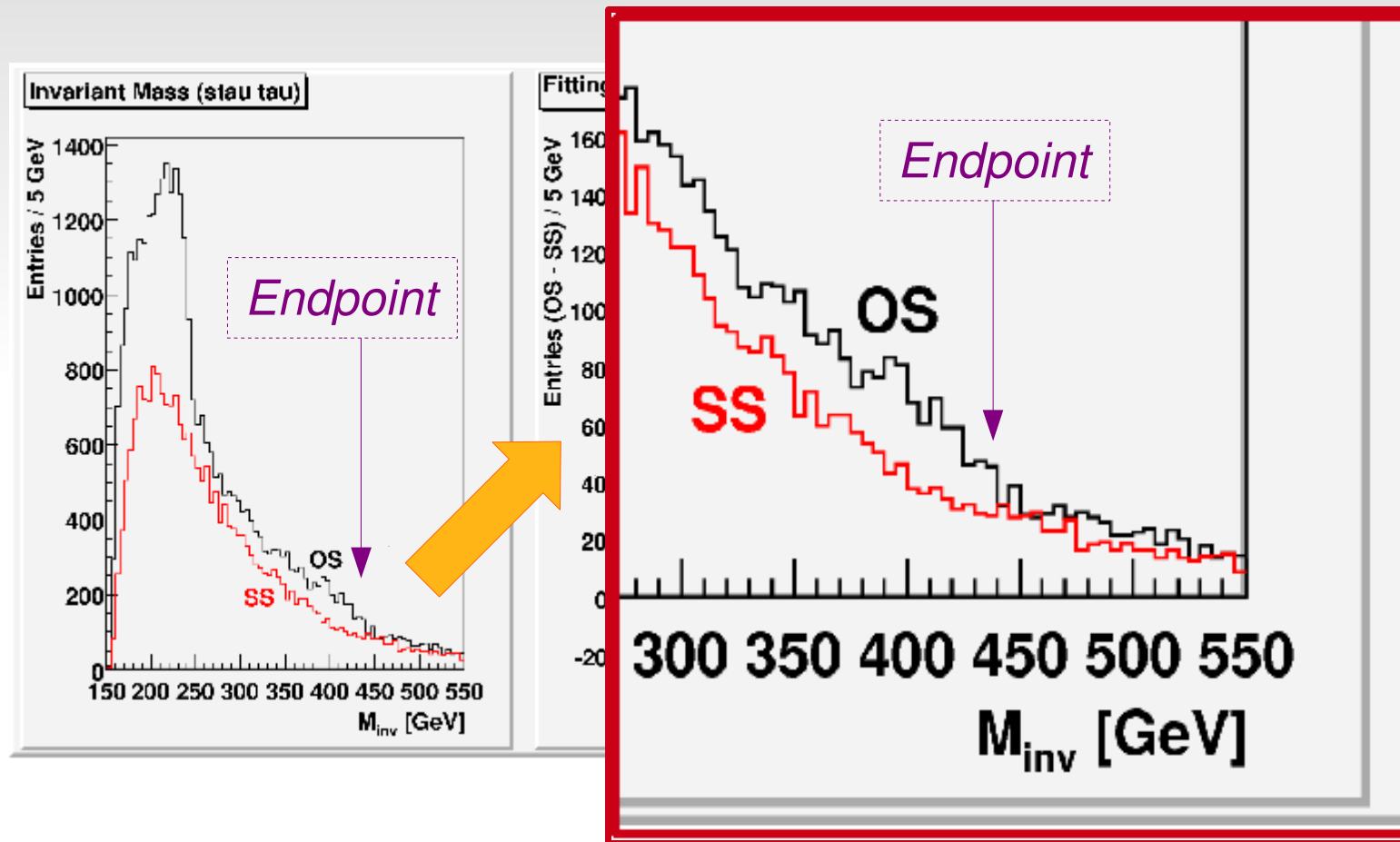
* From the error of stau mass

$$\pm 100 \text{ MeV}$$

→ Bino mass measurement : $\delta m_{\tilde{B}} \sim 1 \text{ GeV}$

Neutralino Masses

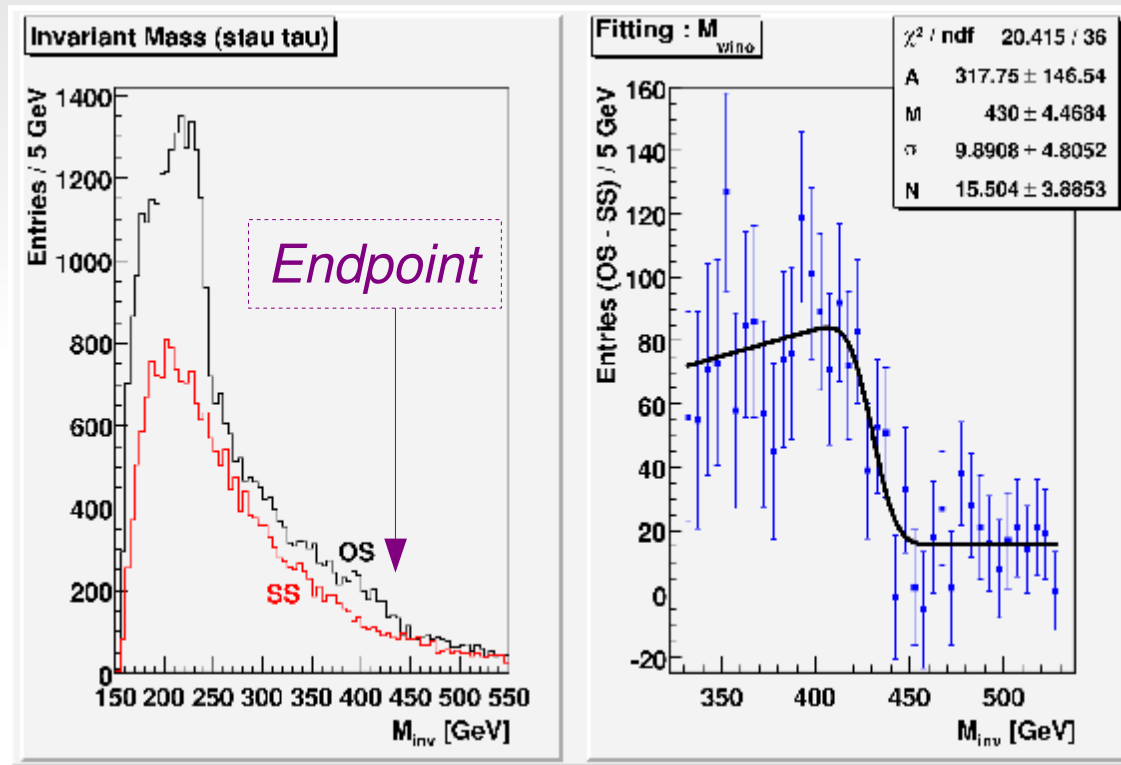
Also, **wino mass** can be measured by using charge subtraction.



This endpoint is *not clear* before we adopt *charge subtraction*.

Neutralino Masses

Also, **wino mass** can be measured by using charge subtraction.



$$M_{\tilde{W}_0} = 430.0 \pm 4.5 \text{ GeV}$$

(425.9 : underlying)

Uncertainties

* Statistics

$$\pm 5 \text{ GeV}$$

* From the error of stau mass

$$\pm 100 \text{ MeV}$$

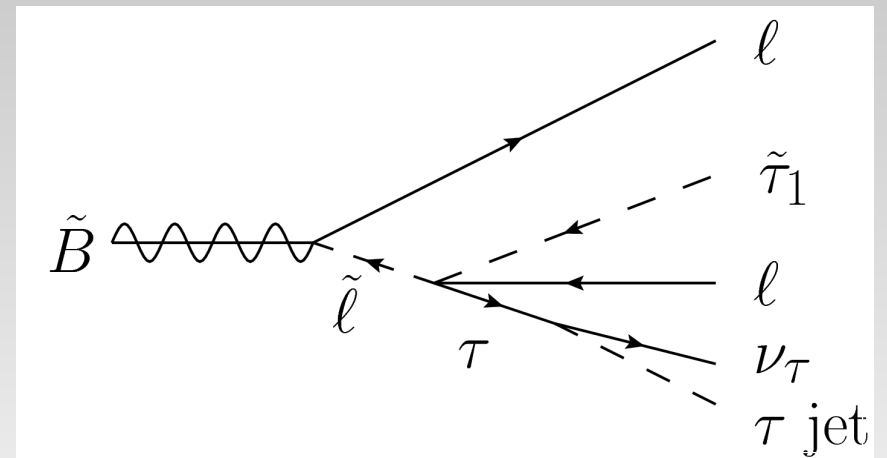
→ Wino mass measurement : $\delta m_{\tilde{W}_0} \sim 5 \text{ GeV}$

Selectron & Smuon Masses

2-step SUSY decay chain:

bino $\rightarrow \ell + \text{slepton}$,

slepton $\rightarrow \ell + \tau + \text{stau}$; followed by $\tau \rightarrow \tau\text{-jet} + \nu_\tau$



Collinear approximation; $p_\tau = r p_{\tau\text{-jet}}$ ($r \geq 1$)

The value of r is determined by the condition:

$$m_{\tilde{B}}^2 = (p_{\ell^+} + p_{\ell^-} + r p_{\tau\text{-jet}} + p_{\tilde{\tau}_1})^2$$

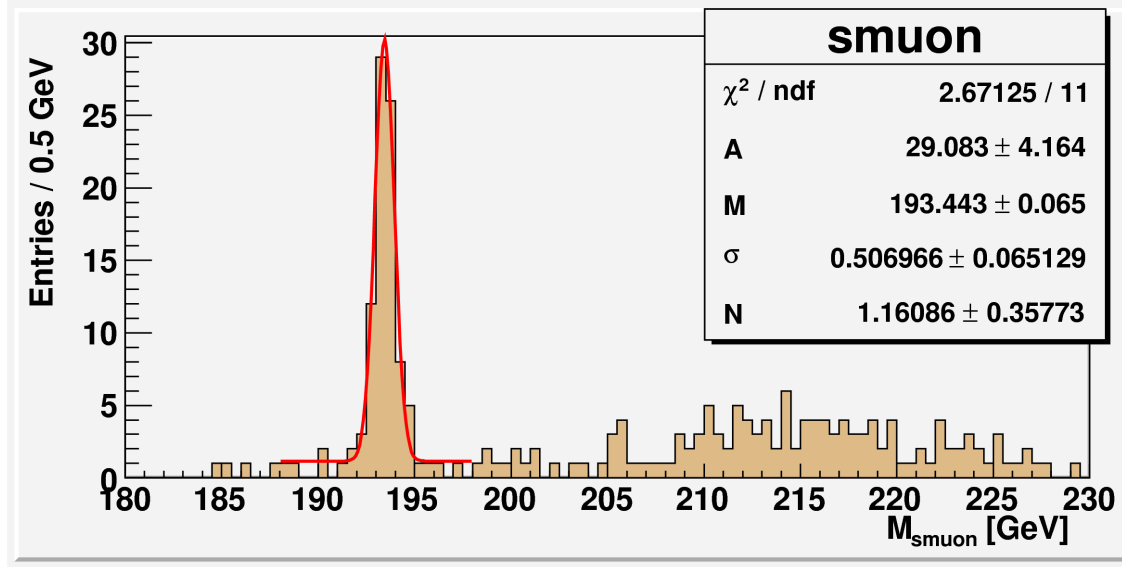
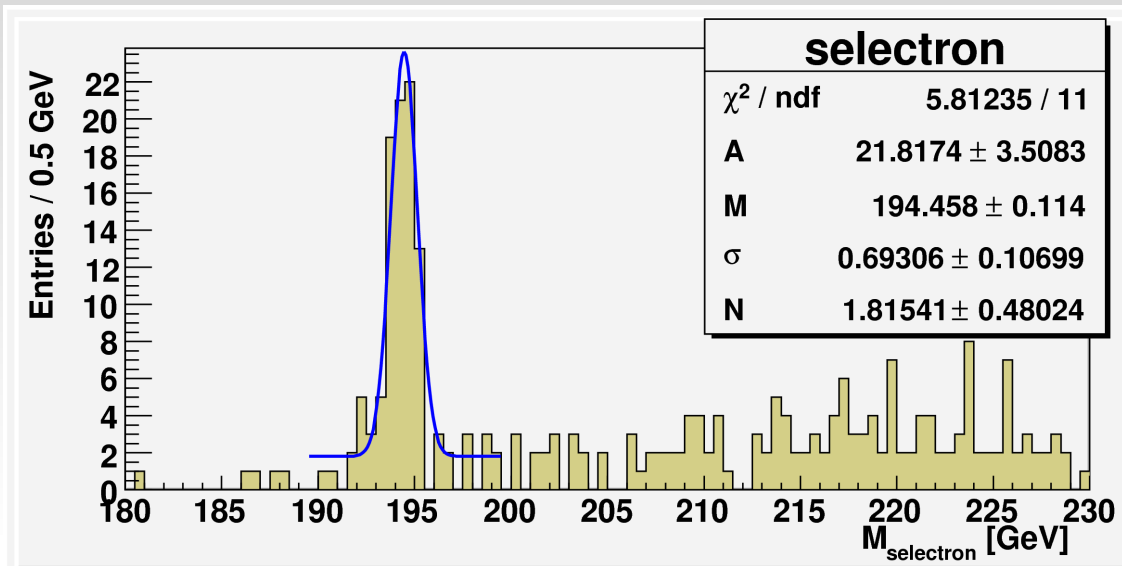
Then, $\longrightarrow M_{\tilde{\ell}} = \sqrt{(p_\ell + r p_{\tau\text{-jet}} + p_{\tilde{\tau}_1})^2}$

Event Selection

(1)(2) At least one pair of stau and τ -tagged jet (OS)

(3) At least one pair of leptons(SF,OS) with $p_T > 15$ GeV

Selectron & Smuon Masses



$$M_{\tilde{e}} = 194.5 \pm 0.1 \text{ GeV}$$

(194.4 : underlying)

$$M_{\tilde{\mu}} = 193.4 \pm 0.1 \text{ GeV}$$

(193.4 : underlying)

Uncertainties

* Statistics

$$\pm 100 \text{ MeV}$$

* From the error of stau mass

$$\pm 100 \text{ MeV}$$

* From the error of bino mass

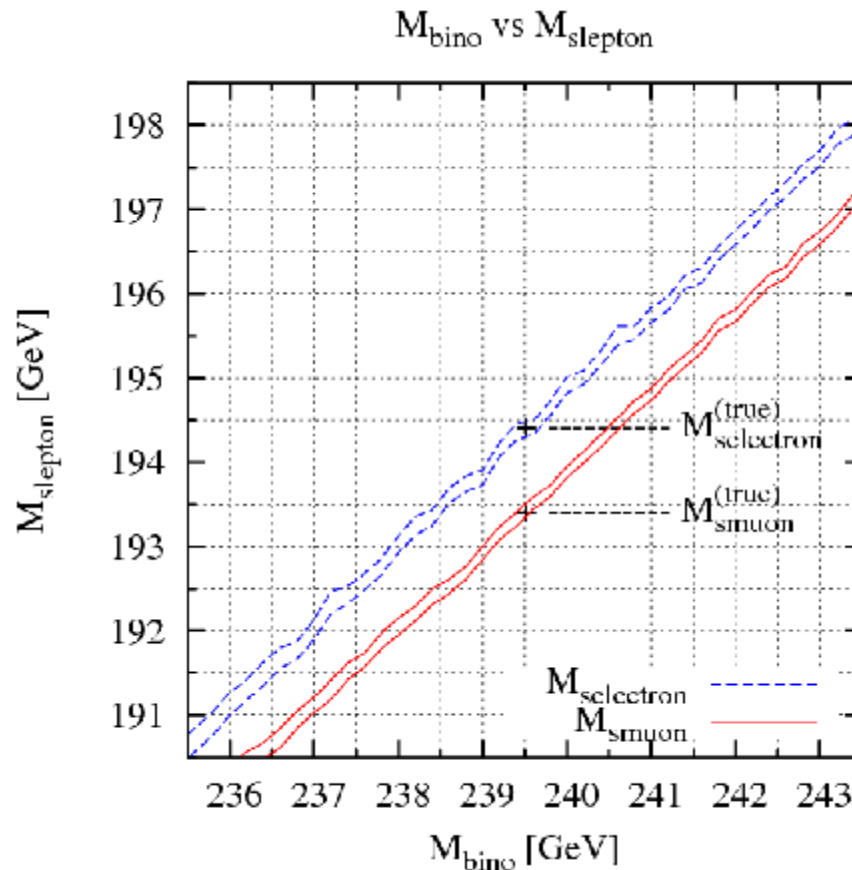
$$\pm 1 \text{ GeV}$$

→ Slepton masses :

$$\delta m_{\tilde{e}} \sim \delta m_{\tilde{\mu}} \sim 1 \text{ GeV}$$

Selectron & Smuon Masses

The mass difference of two sleptons, $M_{\tilde{e}} - M_{\tilde{\mu}}$, is **not sensitive** to the uncertainty of the bino mass.



→ Mass difference :
 $\delta(m_{\tilde{e}} - m_{\tilde{\mu}})$
 $\sim 100 \text{ MeV}$

Good accuracy!

Selectron & Smuon Masses

The mass difference measurements with the good accuracy.

→ **Implications for the SUSY breaking mechanism**

* Loop effects to the slepton masses (Yukawa interaction)

$$\rightarrow M_{\tilde{e}} - M_{\tilde{\mu}} \sim O(100) \text{ MeV} \quad \text{for the large } \tan\beta$$

* The size of SUGRA effects are estimated $\sim m_{3/2}^2/m_{\tilde{\ell}}$

$$\rightarrow \text{if } m_{3/2} \sim \text{a few GeV, } M_{\tilde{e}} - M_{\tilde{\mu}} \sim O(100) \text{ MeV}$$

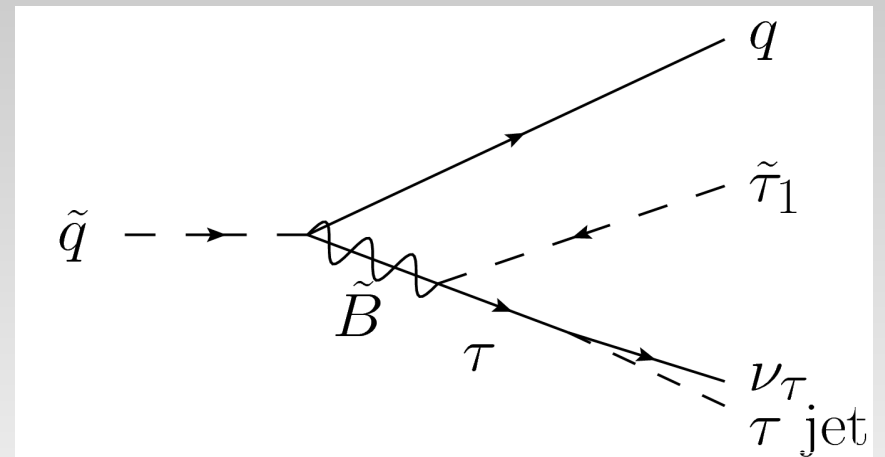
→ *These effects are detectable.*

Squark Masses

2-step SUSY decay chain:

squark \rightarrow q + bino,

bino \rightarrow τ + stau ; followed by $\tau \rightarrow$ τ -jet + ν_τ



τ 's momentum is reconstructed by $m_{\tilde{B}}^2 = (rp_{\tau\text{-jet}} + p_{\tilde{\tau}_1})^2$

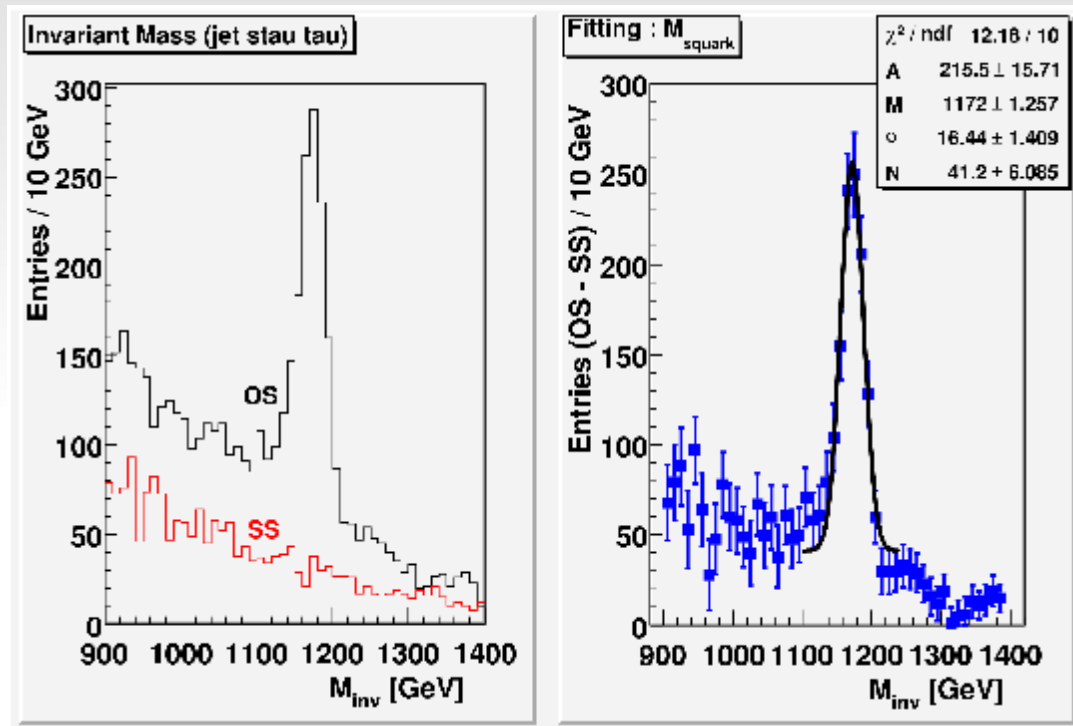
Then, $\longrightarrow M_{\tilde{q}} = \sqrt{(p_{\text{jet}} + rp_{\tau\text{-jet}} + p_{\tilde{\tau}_1})^2}$

Event Selection

- (1)(2) At least one pair of stau and τ -tagged jet
- (3) At least one jet with $p_T > 100$ GeV
- (4) No isolated leptons with $p_T > 15$ GeV

Squark Masses

We use upto leading 4 high- p_T jets, and perform **charge subtraction**.



$$M_{\tilde{q}} = 1172 \pm 1 \text{ GeV} \\ (1180 : \text{Squark(R)})$$

$$[\text{BR}(\tilde{q}_R \rightarrow q\tilde{\chi}_1^0) \simeq 100\%]$$

Uncertainties

- * Statistics + *Systematics*
(± 1 GeV) + (- 8 GeV)
- * From the error of stau mass
± 100 MeV
- * From the error of bino mass
± 1 GeV

→ Squark mass measurement : $\delta m_{\tilde{q}} \sim 10 \text{ GeV}$

Conclusions

We have discussed mass measurements of superparticles in the long-lived stau scenario at the LHC experiments.

- * Neutralino Masses

by endpoint analysis : $\delta m_{\tilde{B}} \sim 1 \text{ GeV}, \delta m_{\tilde{W}_0} \sim 5 \text{ GeV}$

charge subtraction method is useful.

- * Selectron & Smuon Masses ($m_{\tilde{B}} > m_{\tilde{\ell}_R} > m_{\tilde{\tau}_1}$)

by peak analysis: $\delta m_{\tilde{e}} \sim \delta m_{\tilde{\mu}} \sim 1 \text{ GeV}$

- * mass difference : $\delta(m_{\tilde{e}} - m_{\tilde{\mu}}) \sim 100 \text{ MeV}$

- * Squark Masses

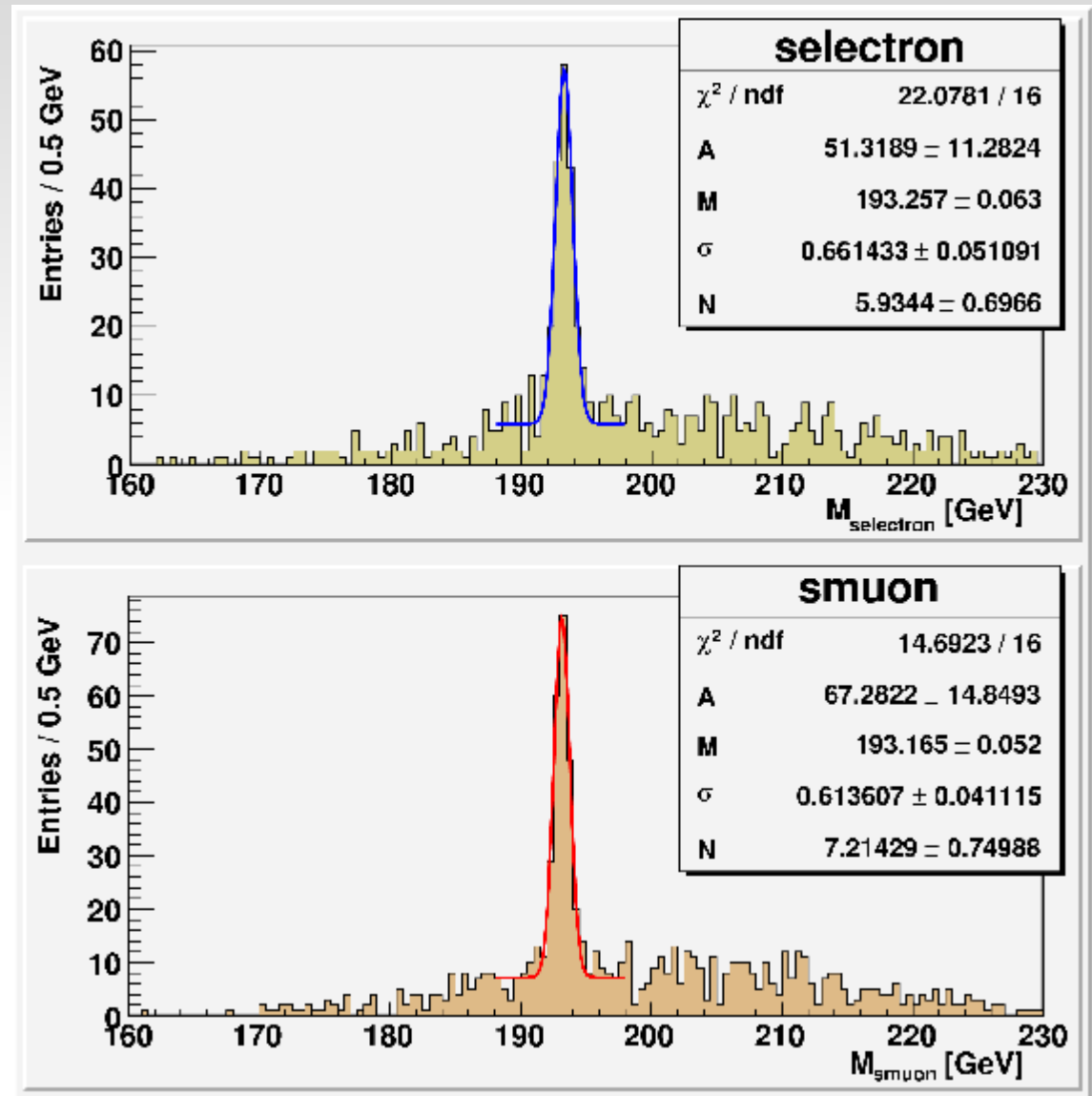
by peak analysis: $\delta m_{\tilde{q}} \sim 10 \text{ GeV}$

informations for the ~~SUSY~~

backup

Another Example; Sweet Spot Model [Ibe,Kitano(07)]

slepton masses

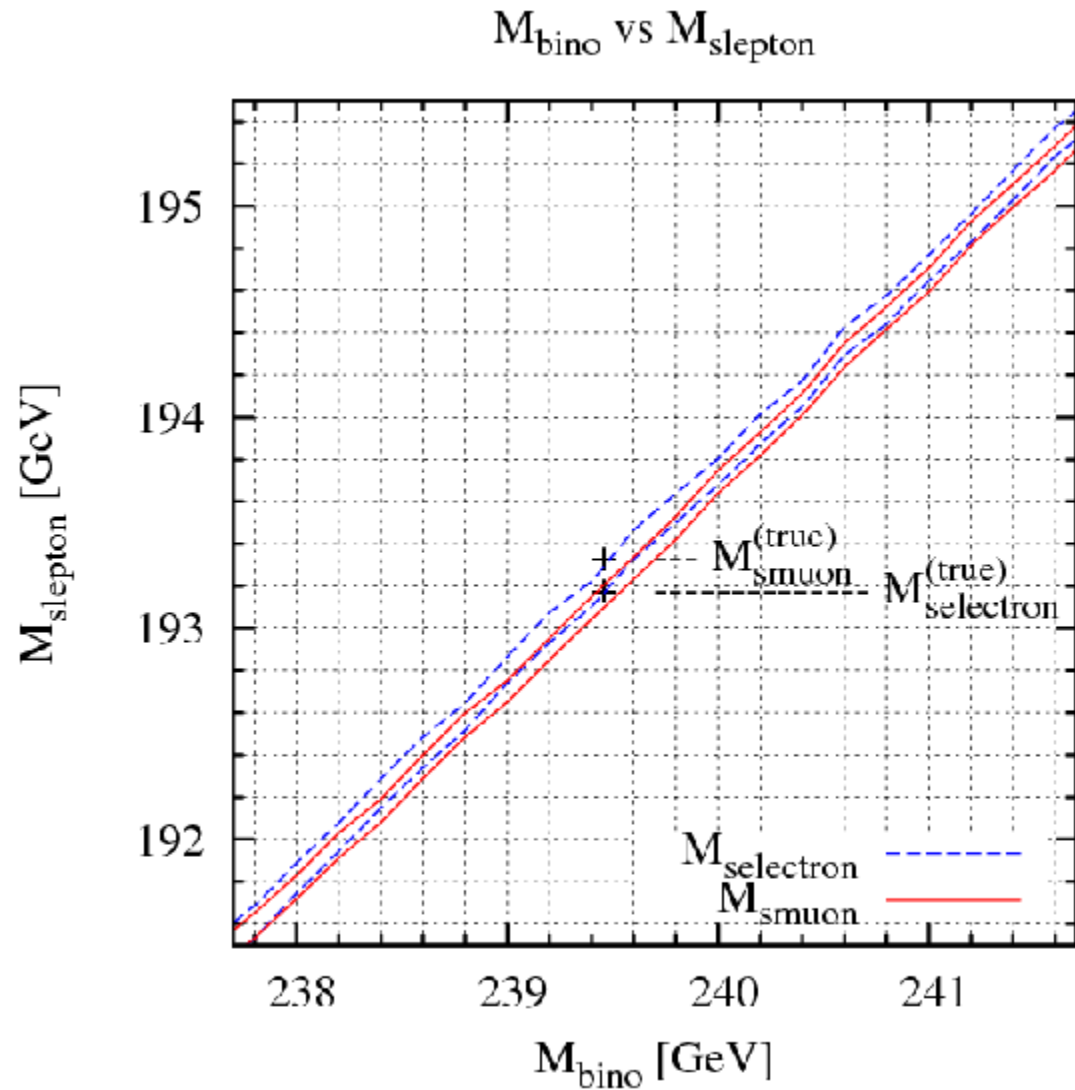


Another Example; Sweet Spot Model [Ibe,Kitano(07)]

bino mass

vs

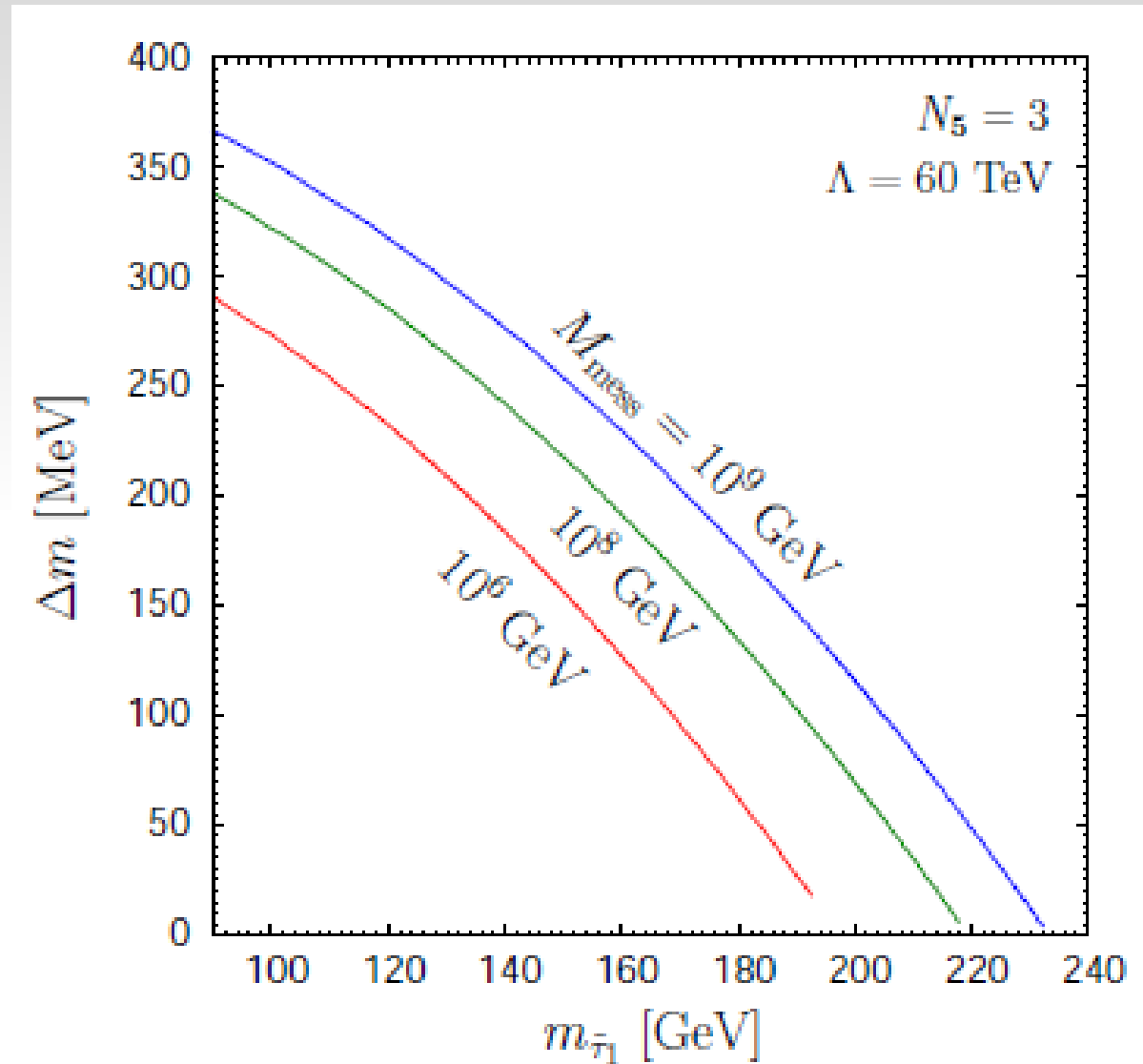
slepton mass



Another Example; Sweet Spot Model [Ibe,Kitano(07)]

*The size of
the mass difference*

*(for the fixed
 N_5 and Λ)*



Introduction

Supersymmetry : A famous extension of the Standard Model

There are several reasons to consider SUSY seriously.

- * Gauge coupling unification at the very high energy scale (*GUT*)
- * Candidates of Dark Matter of the Universe
- * Predicts a light higgs (preferred by EW precision measurements)
- * Naturally solves gauge hierarchy problem

→ It will appear around TeV scale.

↔ in the energy range of the LHC experiments!

So, Question is

What is a signature of SUSY at the LHC?