

# Experimental search for the cosmic background neutrino decay in the cosmic far-infrared background

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for Neutrino Decay Collaboration

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- Summary

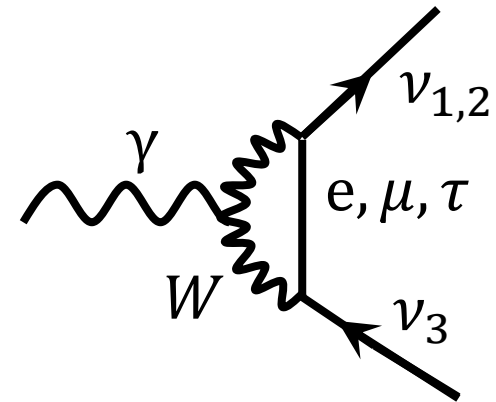
# Neutrino

- Neutrino has 3 mass generations ( $\nu_1, \nu_2, \nu_3$ )
- Neutrino flavor states ( $\nu_e, \nu_\mu, \nu_\tau$ ) are not mass eigenstates

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

→ Neutrino can decay through the loop diagram

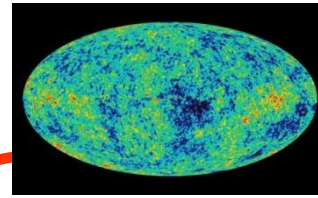
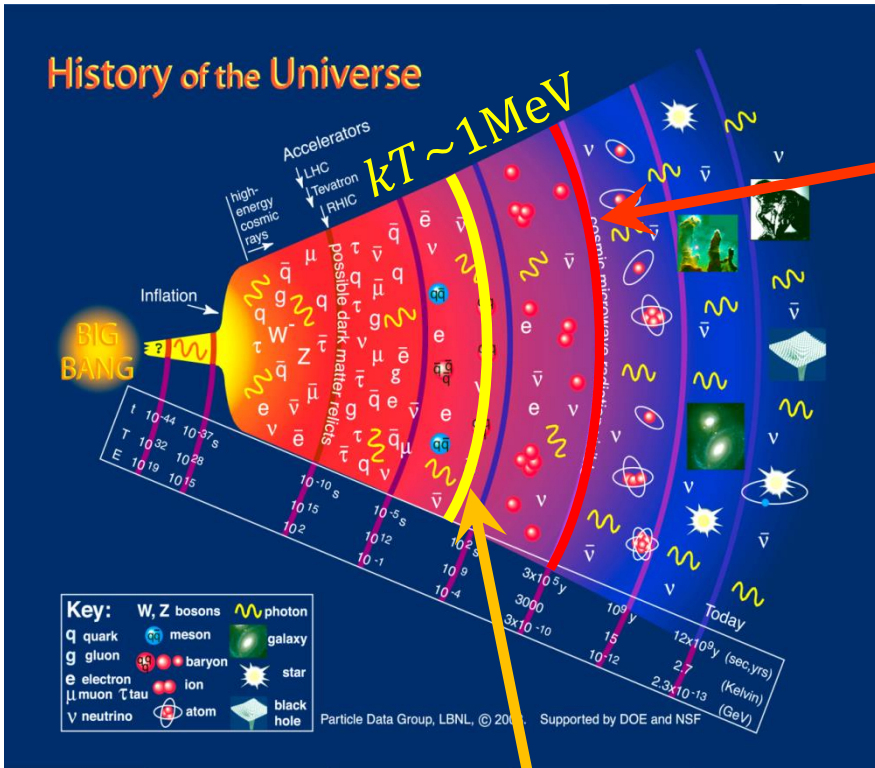
$$- \nu_3 \rightarrow \nu_{1,2} + \gamma$$



- Neutrino lifetime is very long

→ Cosmic neutrino background (CνB) is the best neutrino source for neutrino decay search

# Cosmic neutrino background (CνB)

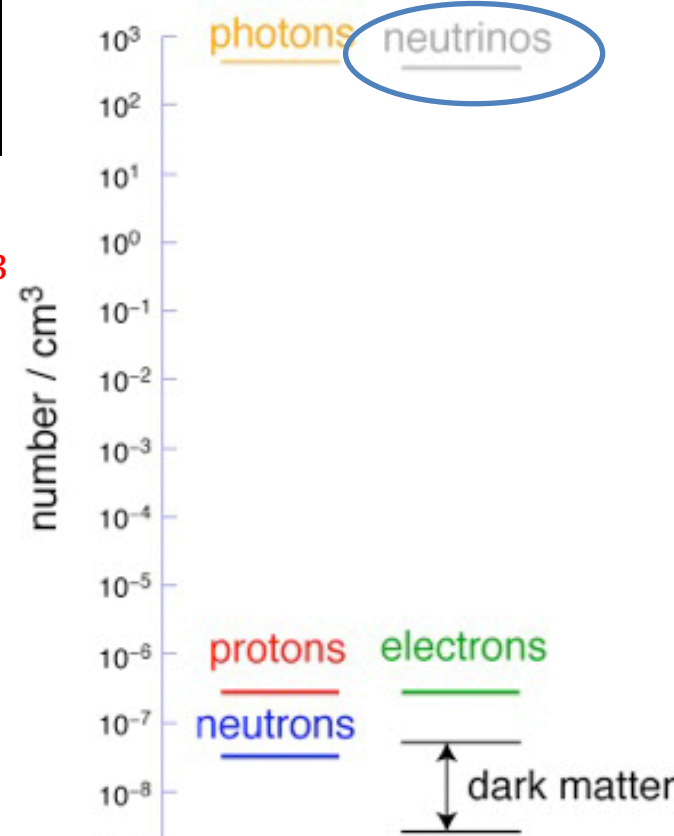


**CMB**

$$n_\gamma = 411/\text{cm}^3$$

$$T_\gamma = 2.73 \text{ K}$$

## The Particle Universe



**CνB (~1s after big bang)**

$$n_\nu + n_{\bar{\nu}} = \frac{3}{4} \left( \frac{T_\nu}{T_\gamma} \right)^3 n_\gamma = 110/\text{cm}^3$$

Not yet  
observed  
experimentally

$$T_\nu = \left( \frac{4}{11} \right)^{\frac{1}{3}} T_\gamma = 1.95 \text{ K}$$

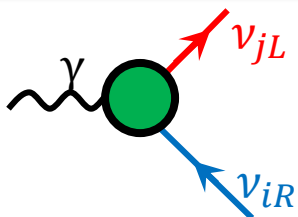
$$\langle p_\nu \rangle = 0.5 \text{ meV}/c$$

# Motivation of $\nu$ -decay search in $C\nu B$

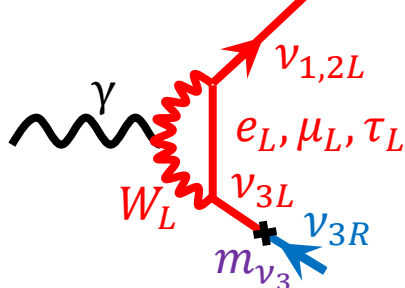
- Search for  $\nu_3 \rightarrow \nu_{1,2} + \gamma$  in cosmic neutrino background ( $C\nu B$ )
  - Direct detection of  $C\nu B$
  - Direct detection of transition magnetic dipole moment of neutrino
  - Direct measurement of neutrino mass:  $m_3 = (m_3^2 - m_{1,2}^2)/2E_\gamma$
- Aiming at sensitivity of detecting  $\gamma$  from  $\nu$  decay for  $\tau(\nu_3) = O(10^{17} \text{ yrs})$ 
  - SM expectation  $\tau = O(10^{43} \text{ yrs})$
  - Current experimental lower limit  $\tau > O(10^{12} \text{ yrs})$
  - L-R symmetric model (for Dirac neutrino) predicts down to  $\tau = O(10^{17} \text{ yrs})$  for  $W_L$ - $W_R$  mixing angle  $\zeta < 0.02$

Neutrino magnetic moment term

$$\bar{\nu}_{jL} \sigma_{\mu\nu} \nu_{iR}$$



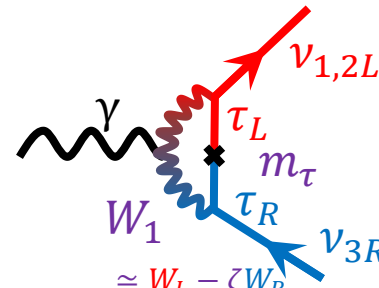
**SM:**  $SU(2)_L \times U(1)_Y$



$$\Gamma \sim (10^{43} \text{ yr})^{-1}$$

Suppressed by  $m_\nu$ , GIM

**LRS:**  $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$



$$\Gamma \sim (10^{17} \text{ yr})^{-1}$$

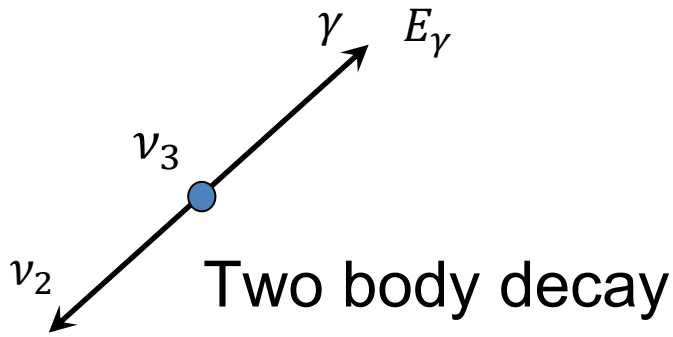
Only suppressed by  $\zeta \sim 0.02$

PRL 38,(1977)1252, PRD 17(1978)1395

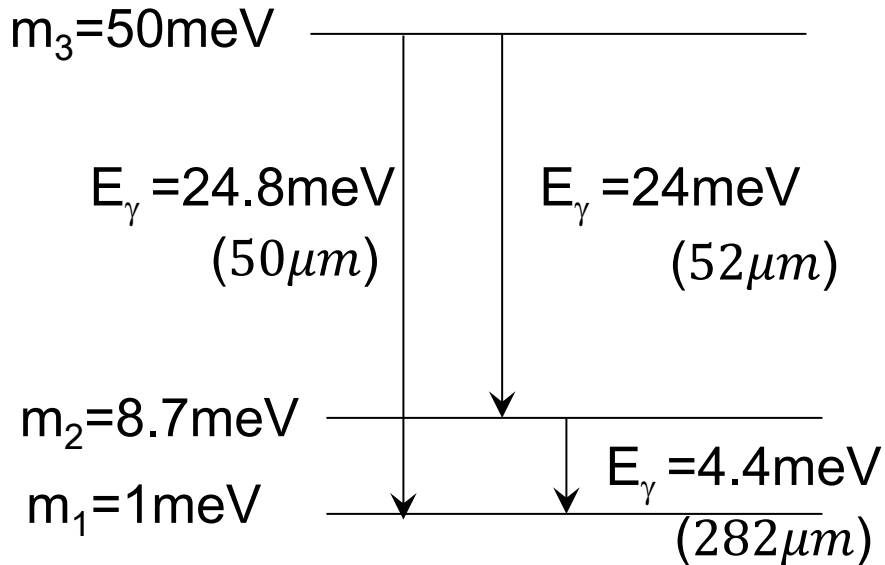
$$\begin{pmatrix} W_1 \\ W_2 \end{pmatrix} = \begin{pmatrix} \cos\zeta & -\sin\zeta \\ \sin\zeta & \cos\zeta \end{pmatrix} \begin{pmatrix} W_L \\ W_R \end{pmatrix}$$

**$10^{26}$   
enhancement to  
SM**

# Photon Energy in Neutrino Decay



$$\nu_3 \rightarrow \nu_{1,2} + \gamma \quad E_\gamma = \frac{m_3^2 - m_{1,2}^2}{2m_3}$$

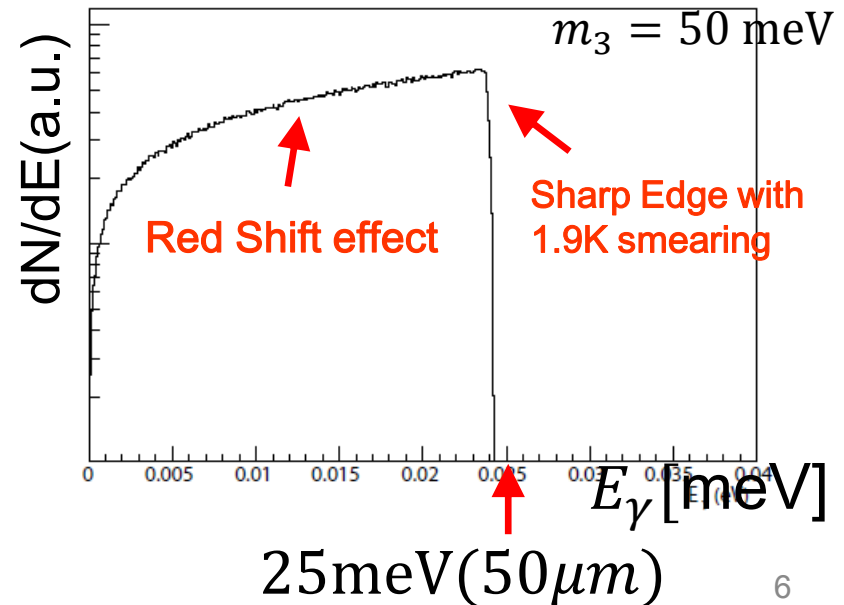


- From neutrino oscillation
  - $\Delta m_{23}^2 = |m_3^2 - m_2^2| = 2.4 \times 10^{-3} \text{ eV}^2$
  - $\Delta m_{12}^2 = 7.65 \times 10^{-5} \text{ eV}^2$
- From Planck+WP+highL+BAO
  - $\sum m_i < 0.23 \text{ eV}$

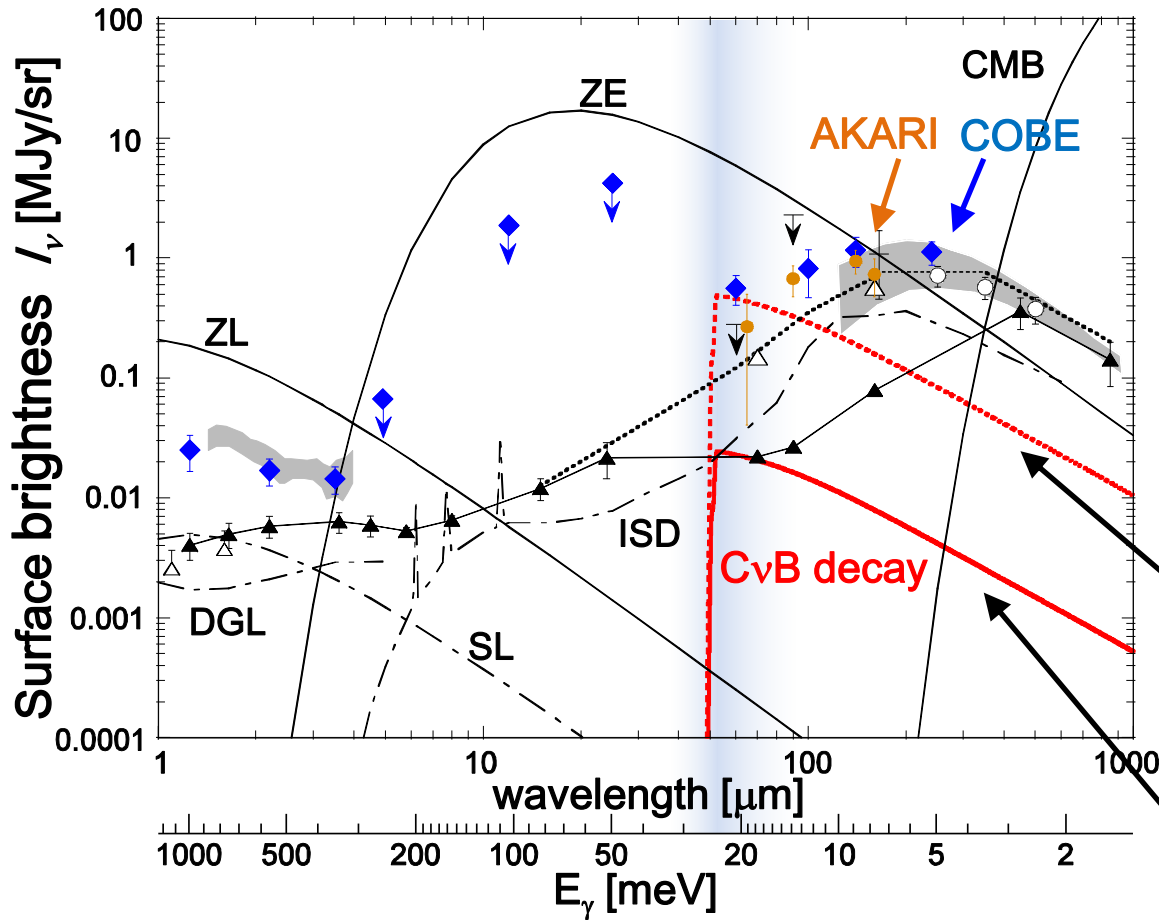
→  $50 \text{ meV} < m_3 < 87 \text{ meV}$

$E_\gamma = 14 \sim 24 \text{ meV}$  ( $\lambda_\gamma = 51 \sim 89 \mu\text{m}$ )

$E_\gamma$  distribution in  $\nu_3 \rightarrow \nu_2 + \gamma$



# Backgrounds to CνB decay



Zodiacal Emission

$$I_\nu \sim 8 \text{ MJy/sr}$$

CIB

$$\lambda I_\lambda \sim 0.1 - 0.5 \text{ MJy/sr}$$

CνB decay

$$\tau = 5 \times 10^{12} \text{ yrs}$$

$$I_\nu \sim 0.5 \text{ MJy/sr}$$

$$\tau = 1 \times 10^{14} \text{ yrs}$$

$$I_\nu \sim 25 \text{ kJy/sr}$$

at  $\lambda = 50 \mu\text{m}$

Expected  $E_\gamma$  spectrum

$$m_3 = 50 \text{ meV}$$

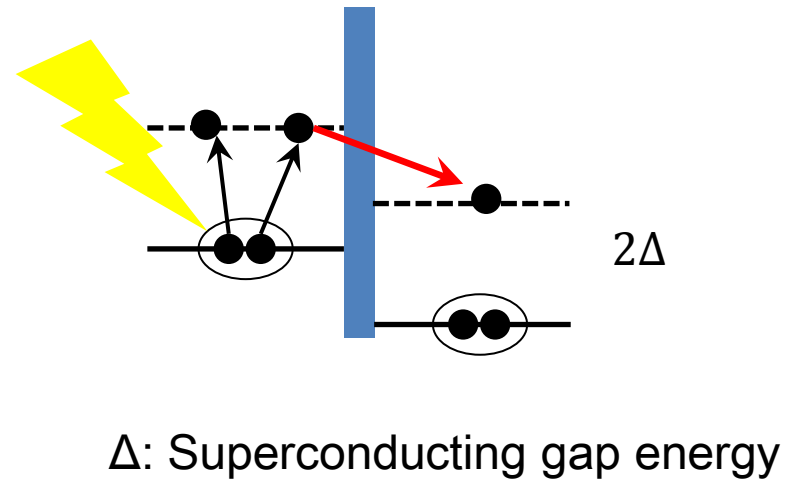
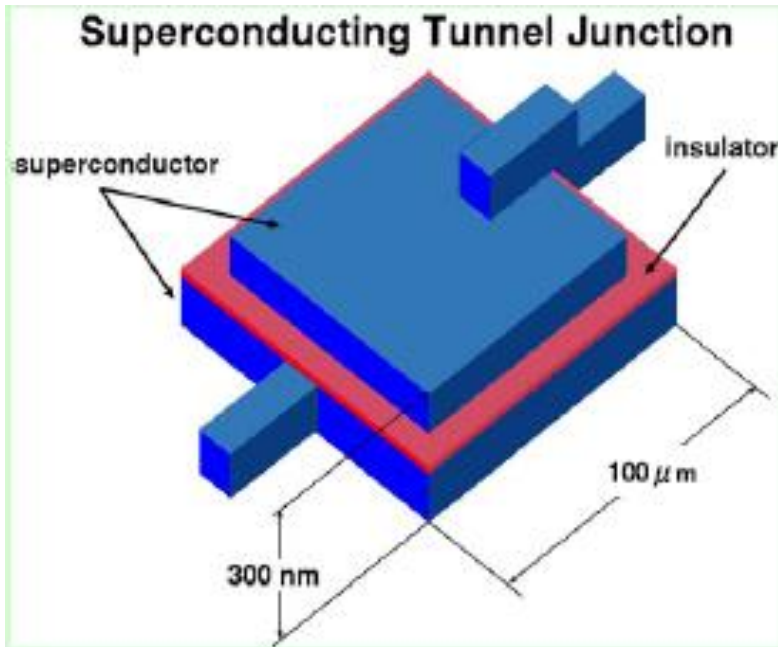
# Detector requirements

- Need continuous spectrum of photon energy around  $\lambda=50\mu\text{m}$  (far infrared photon) with highly precise accuracy
  - **Photon-by-photon energy measurement** with better than 2% resolution for  $E_\gamma = 25\text{meV}$  ( $\lambda=50\mu\text{m}$ ) to achieve better S/N as well as to identify the sharp edge in the spectrum
  - A ground based experiment is impossible, so rocket and/or satellite experiment with this detector is required
- **Superconducting Tunneling Junction (STJ) detectors in development**
  - Array of 50 Nb/Al-STJ pixels with diffraction grating covering  $\lambda = 40 - 80\mu\text{m}$ 
    - **For rocket experiment aiming at improvement of current lower limit for  $\tau(\nu_3)$  by 2 order :  $O(10^{14}\text{yr})$  in a 200-sec measurement**
  - STJ using Hafnium: Hf-STJ for satellite experiment
    - $\Delta = 20\mu\text{eV}$  : Superconducting gap energy for Hafnium
    - $N_{\text{q.p.}} = 25\text{meV}/1.7\Delta = 735$  for 25meV photon:  $\Delta E/E < 2\%$  if Fano-factor is less than 0.3



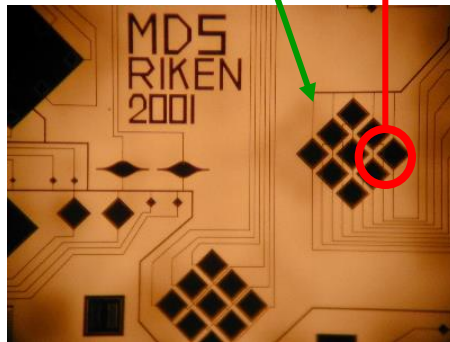
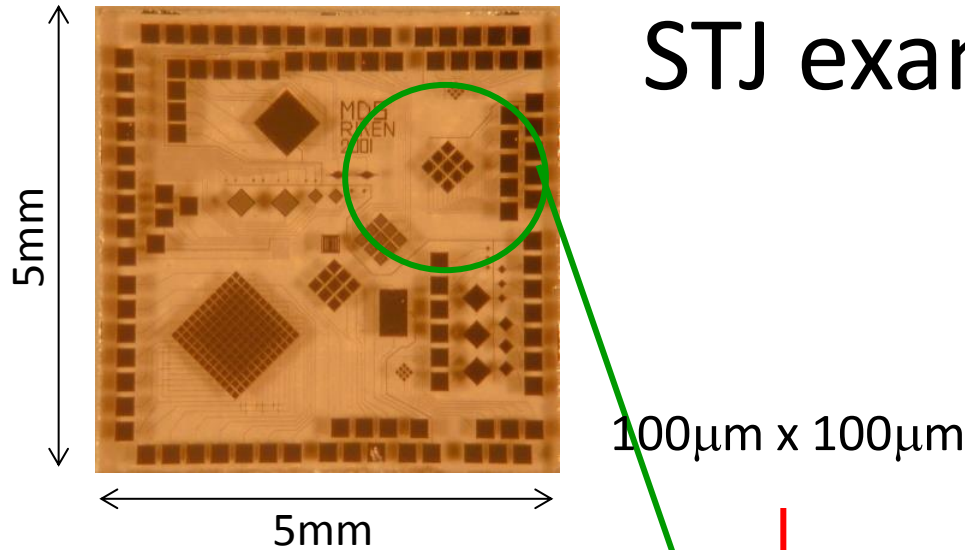
# STJ(Superconducting Tunnel Junction ) Detector

- Superconducting / **Insulator** /Superconducting Josephson junction device

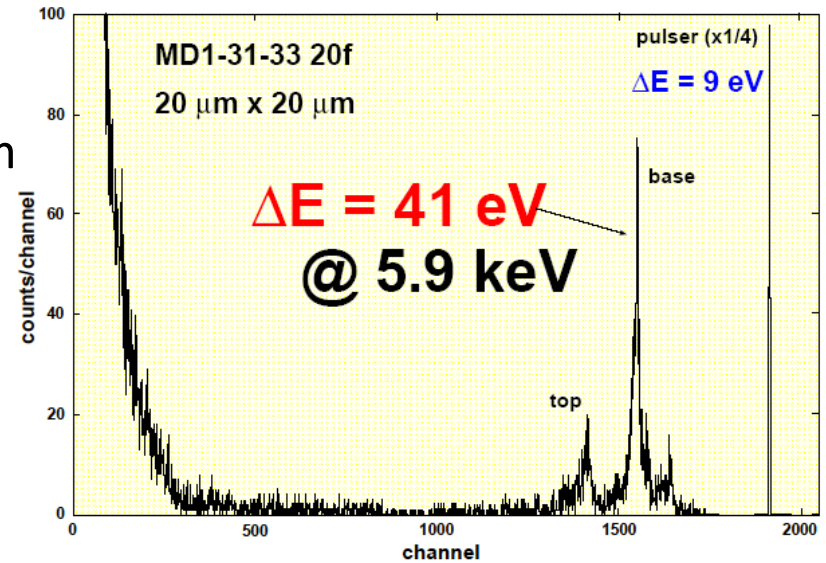


A bias voltage ( $|V| < 2\Delta$ ) is applied across the junction.  
A photon absorbed in the superconductor breaks Cooper pairs and creates tunneling current of quasi-particles proportional to the photon energy.

# STJ examples



5.9KeV X-ray



H. Sato (RIKEN)

- STJs are already in practical use as a single photon spectrometer for a photon ranging from near-infrared to X-ray, and shows excellent performances comparing to conventional semiconductor detectors

**But no example for far-infrared photon**

# STJ energy resolution

Statistical fluctuation in number of quasi-particles determines STJ energy resolution

➔ Smaller superconducting gap energy  $\Delta$  yields better energy resolution

$$\sigma_E = \sqrt{(1.7\Delta)FE}$$

$\Delta$ : Superconducting gap energy  
 F: fano factor  
 E: Photon energy

	Si	Nb	Al	Hf
T <sub>c</sub> [K]		9.23	1.20	0.165
$\Delta$ [meV]	1100	1.550	0.172	0.020

T<sub>c</sub> :SC critical temperature  
 Need  $\sim 1/10T_c$  for practical operation

## Nb

Well-established as Nb/Al-STJ (back-tunneling gain from Al-layer)

$$N_{q.p.} = 25\text{meV}/1.7\Delta = 9.5$$

Poor energy resolution, but photon counting is possible in principle

## Hf

Hf-STJ as a photon detector is not established

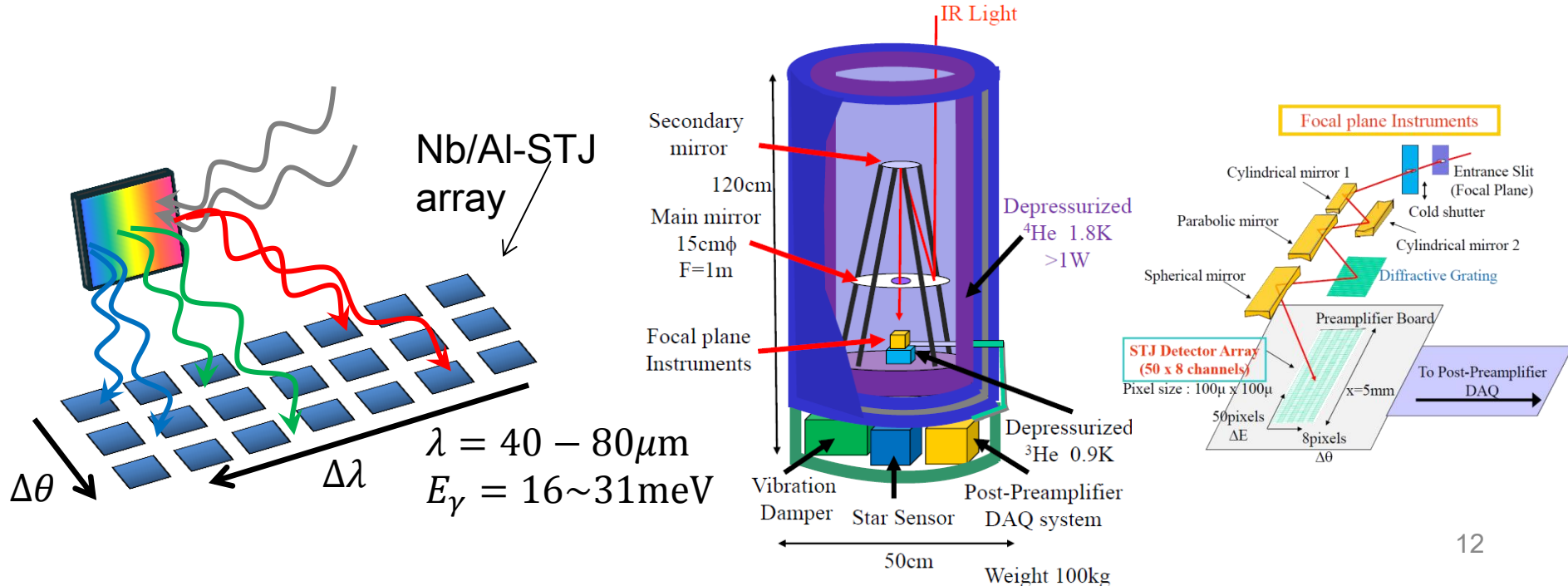
$$N_{q.p.} = 25\text{meV}/1.7\Delta = 735$$

2% energy resolution is achievable if Fano factor  $< 0.3$

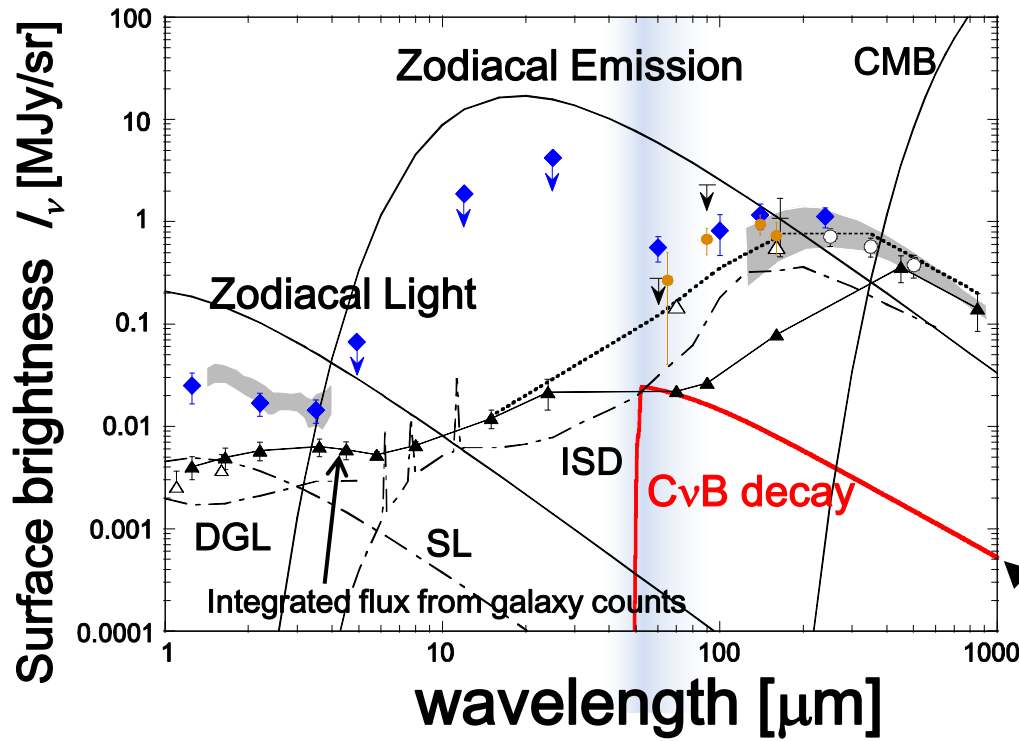
**In both cases, developments are challenging**

# Proposal of a rocket experiment

- Expect 200s measurement at altitude of 200~300km
  - Telescope with diameter of 15cm and focal length of 1m
  - All optics (mirrors, filters, shutters and grating) will be cooled below 4K
- Diffraction grating covering  $\lambda=40-80\mu\text{m}$  (16-31meV) and array of Nb/Al-STJ pixels:  $50(\lambda)\times 8(\theta)$ 
  - Use each Nb/Al-STJ pixel as **a single-photon counting detector** for FIR photon of  $\lambda=40-80\mu\text{m}$  ( $\Delta\lambda = 0.8\mu\text{m}$ )
  - sensitive area of  $100\mu\text{m}\times 100\mu\text{m}$  for each pixel ( $100\mu\text{rad} \times 100\mu\text{rad}$ )



# Expected accuracy in the spectrum measurement



## Telescope parameters

- Main mirror
  - D=15cm, F=1m
- detector
  - sensitive area  
100 $\mu$ m $\times$ 100 $\mu$ m/ pixel
  - 50x8 array

$$\Delta\lambda = \frac{40\mu\text{m}}{50} = 0.8\mu\text{m}$$

$$\tau = 1 \times 10^{14} \text{yr}$$

- Zodiacal emission  $\Rightarrow$  **343Hz / pixel**
  - 200sec measurement: 0.55M events / 8 pixels (at  $\lambda = 50\mu\text{m}$ )
  - 0.13% accuracy measurement for each wavelength:  **$\delta(I_\nu)=11\text{kJy/sr}$**
- Neutrino decay ( $m_3 = 50 \text{ meV}$ ,  $\tau_\nu = 1 \times 10^{14} \text{yr}$ ):  **$I_\nu=25\text{kJy/sr}$** 
  - **2.3 $\sigma$**  away from statistical fluctuation in ZE measurement

**$\nu$  decay with  $\tau_\nu = 10^{14}$  yrs is possible to detect, or set lower limit!**

# Summary

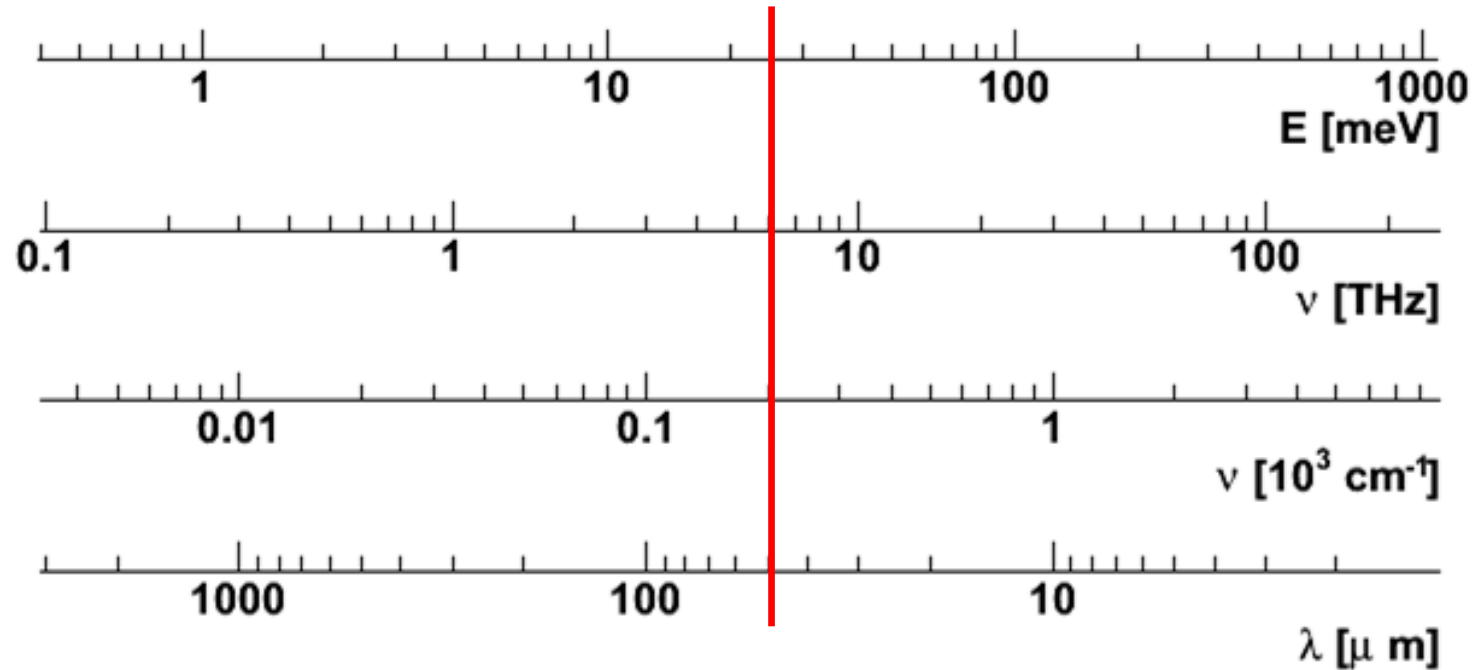
- We propose an experiment to search for neutrino radiative decay in cosmic neutrino background.
- Requirements for the detector is an ability of **photon-by-photon energy measurement with better than 2% energy resolution for  $E_\gamma = 25 \text{ meV}$  ( $\lambda = 50\mu\text{m}$ )**
- Nb/Al-STJ array with grating and Hf-STJ are considered for the experiments and under development.

**Status of Nb/Al-STJ array development will be presented in T. Okudaira's talk**

- It is possible to improve the neutrino lifetime lower limit up to  $O(10^{14}\text{yrs})$  for 200-sec measurement in a rocket experiment with the detector.

# Backup

# Energy/Wavelength/Frequency



$$E_\gamma = 25 \text{ meV}$$

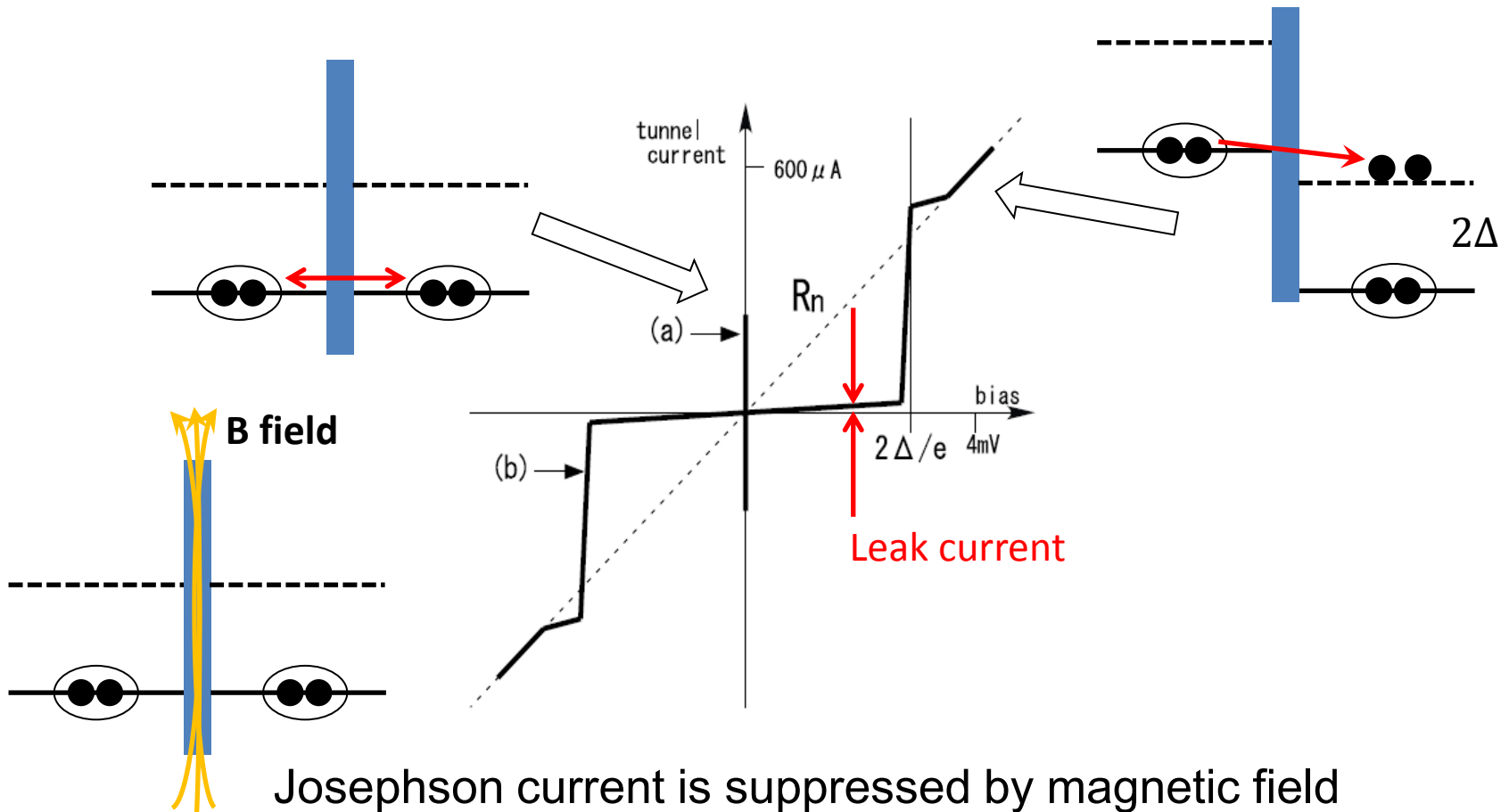
$$\nu = 6 \text{ THz}$$

$$\lambda = 50 \mu\text{m}$$



# STJ I-V curve

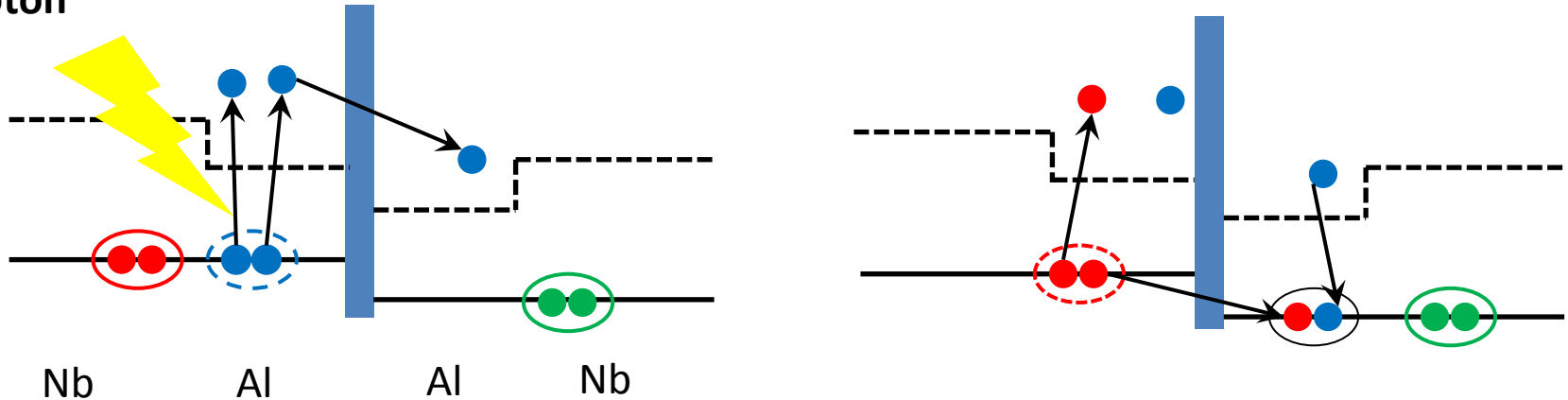
- Sketch of a current-voltage (I-V) curve for STJ
- ➔ The Cooper pair tunneling current (DC Josephson current) is seen at  $V = 0$ , and the quasi-particle tunneling current is seen for  $|V| > 2\Delta$



# STJ back tunneling effect

- Quasi-particles near the barrier can mediate Cooper pairs, resulting in true signal gain
  - Bi-layer fabricated with superconductors of different gaps  $\Delta_{\text{Nb}} > \Delta_{\text{Al}}$  to enhance quasi-particle density near the barrier
  - Nb/Al-STJ Nb(200nm)/Al(10nm)/AlOx/Al(10nm)/Nb(100nm)
- Gain: 2 ~ 200

Photon



## Feasibility of VIS/NIR single photon detection

- Assume typical time constant from STJ response to pulsed light is  $\sim 1\mu s$
- Assume leakage is 160nA

$$160nA = e \times 10^{12}/s = e \times 10^6/\mu s$$

Fluctuation from electron statistics in  $1\mu s$  is

$$e \times \sqrt{10^6}/\mu s = 10^3 e/\mu s$$

While expected signal for 1eV are

(Assume back tunneling gain x10)

$$1eV/1.7\Delta \times 10e = \frac{1eV}{1.7 \times 1.5meV} \times 10 = 4 \times 10^3 e$$

More than 3sigma away from leakage fluctuation

# Feasibility of FIR single photon detection

- Assume typical time constant from STJ response to pulsed light is  $\sim 1\mu s$
- Assume leak current is  $0.1nA$

$$0.1nA = 6.25 \times 10^8 e/s = 6.25 \times 10^2 e/\mu s$$

Fluctuation due to electron statistics in  $1\mu s$  is

$$\sqrt{6.25 \times 10^2 e/\mu s} = 25 e/\mu s$$

While expected signal charge for  $25meV$  are

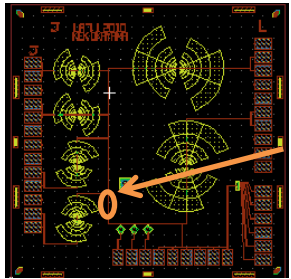
$$25meV/1.7\Delta \times 10e = \frac{25meV}{1.7 \times 1.5meV} \times 10e = 98e$$

(Assume back tunneling gain x10)

More than 3sigma away from leakage fluctuation

- Requirement for amplifier
  - Noise  $< 16e$
  - Gain:  $1V/fC \rightarrow V=16mV$

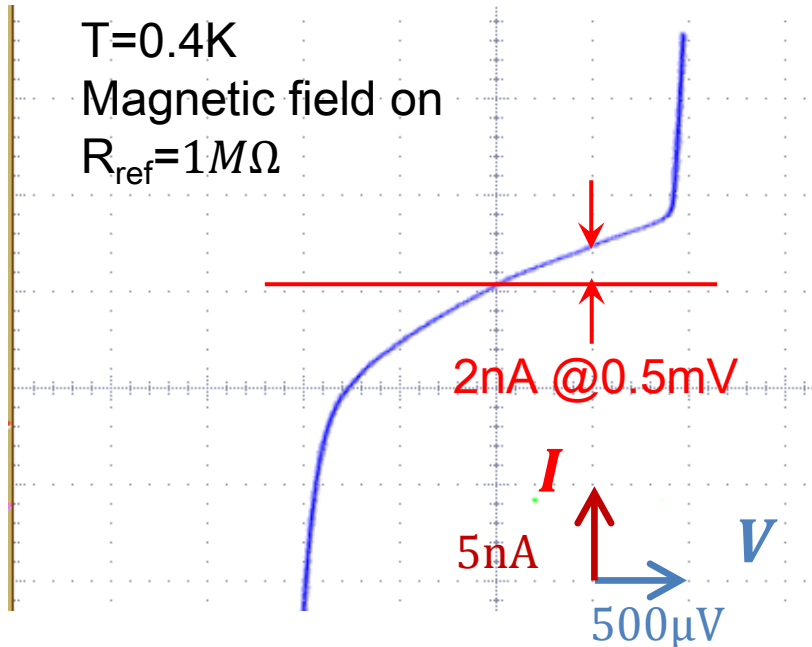
# Temperature dependence of Nb/Al-STJ leak current



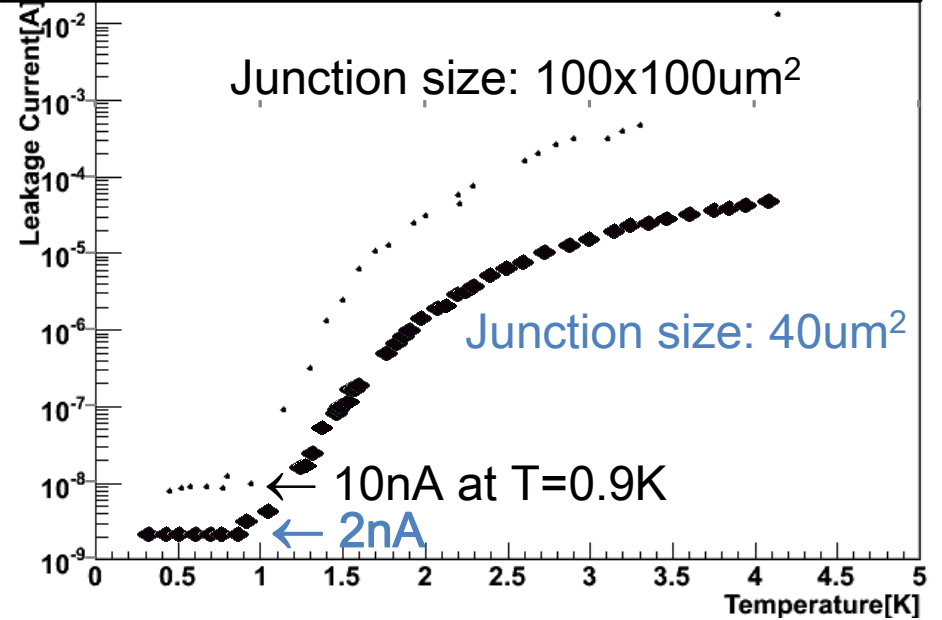
This Nb/Al-STJ is produced by S. Mima (Riken)

Need  $I_{\text{leak}} < 0.1 \text{ nA}$  for single photon counting with  $S/N > 10^3$

40 $\mu\text{m}^2$  Nb/Al-STJ I-V curve



Temperature dependence of leak current



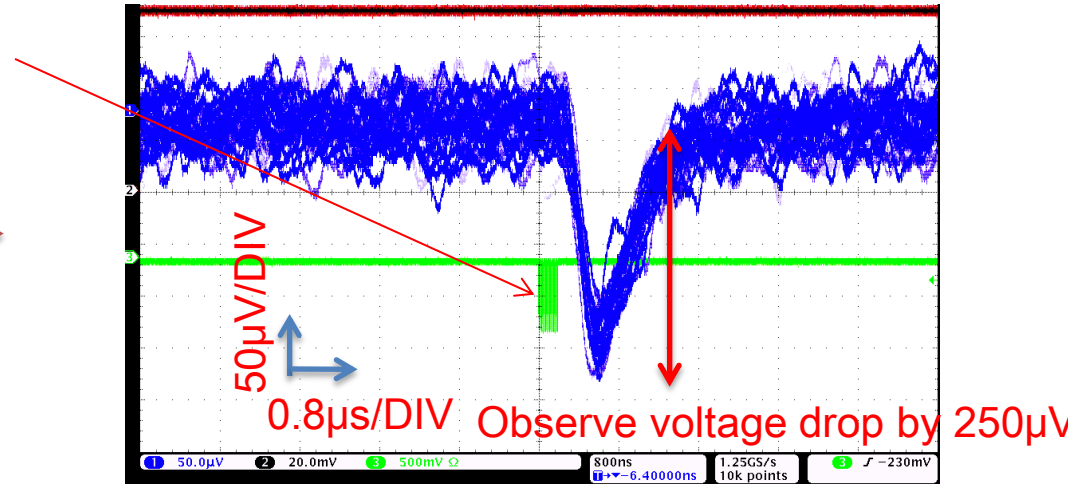
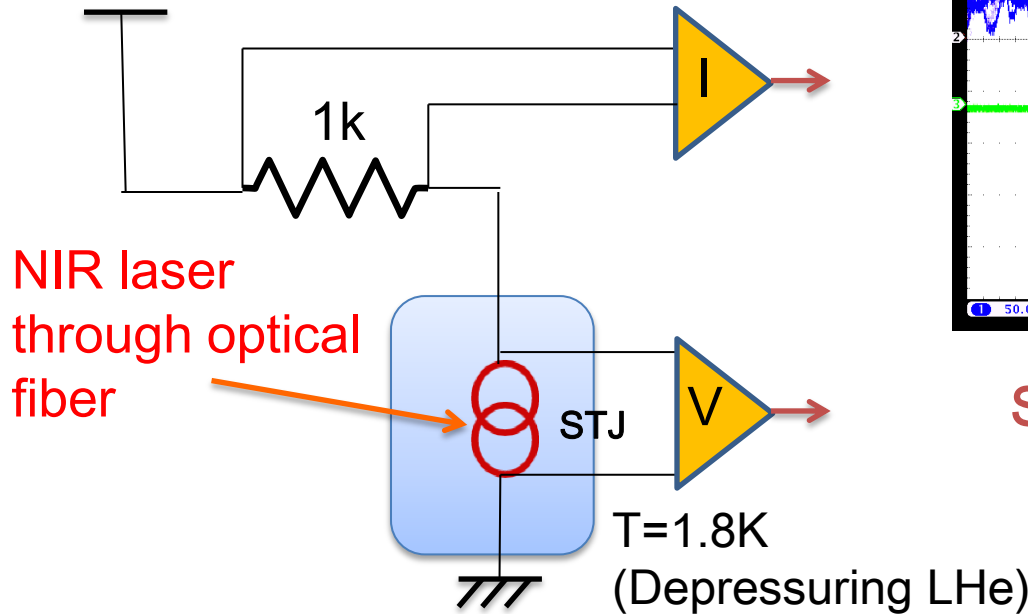
Leak current can be reduced by using small junction size. We are testing STJ of 4 $\mu\text{m}^2$  junction size

Need  $T < 0.9\text{K}$  for detector operation  
 → Use  $^3\text{He}$  sorption or ADR for the operation in rocket experiment

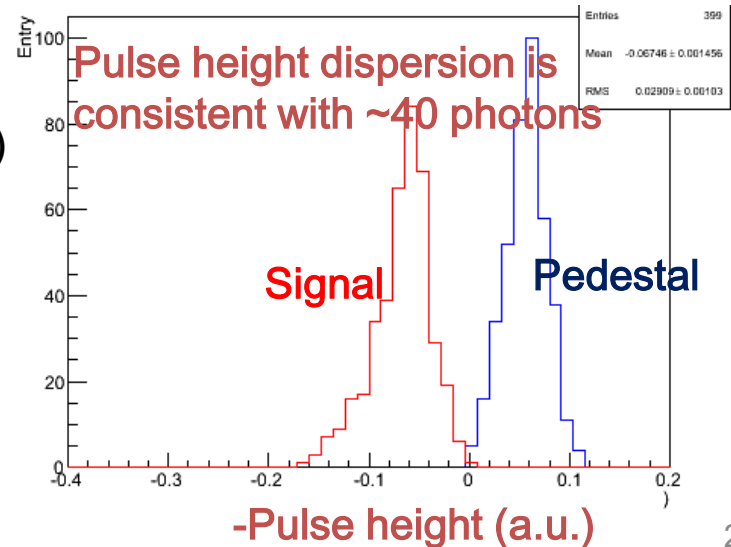
# 100x100 $\mu\text{m}^2$ Nb/Al-STJ response to NIR multi-photons

## Response to NIR laser pulse ( $\lambda=1.31\mu\text{m}$ )

10 laser pulses in 200ns



## Signal pulse height distribution



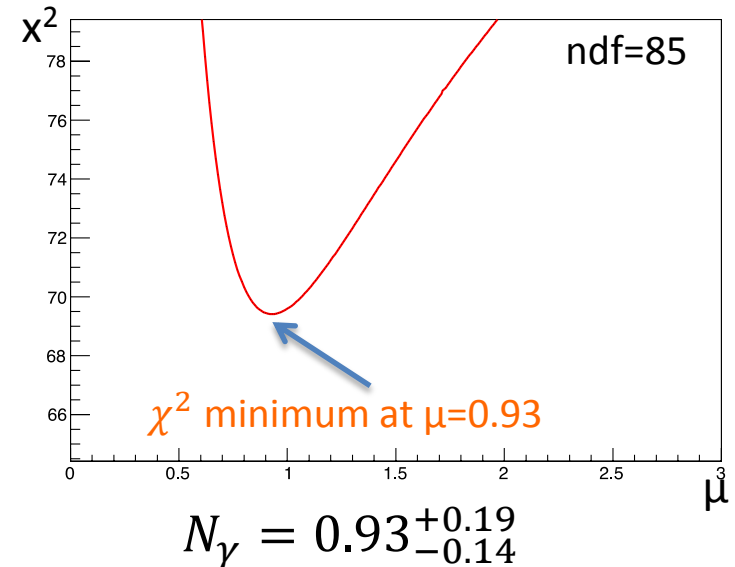
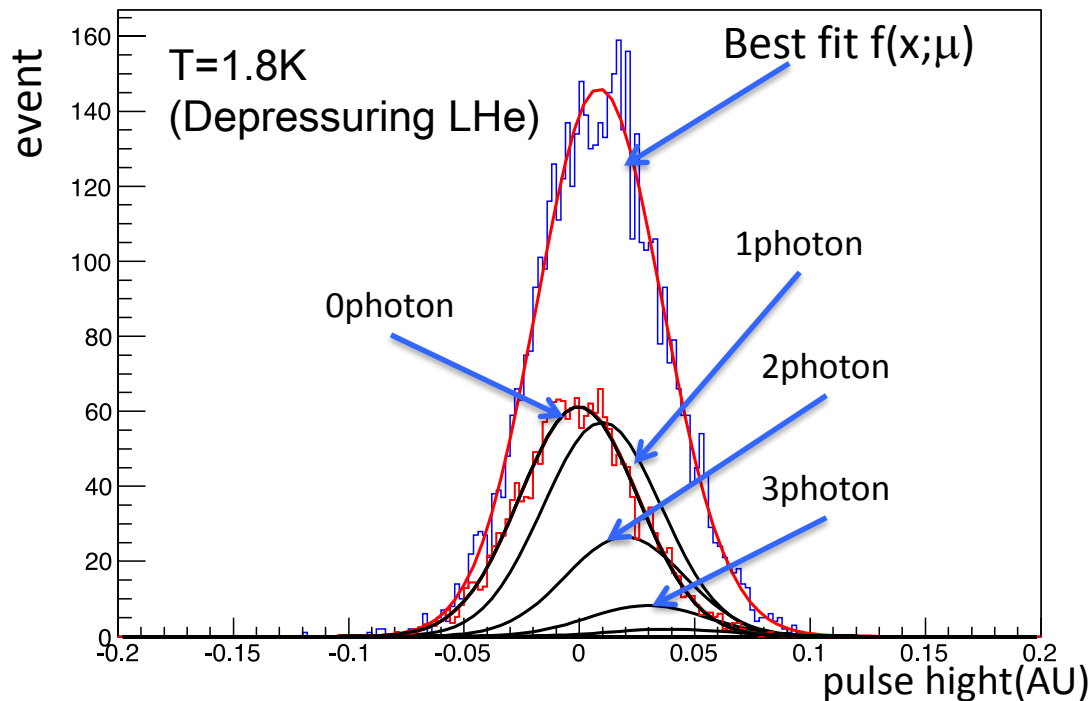
We observed a response to NIR photons

- Response time  $\sim 1\mu\text{s}$
- Corresponding to 40 photon detection (estimated by statistical fluctuations in number of detected photons)

# 4 $\mu\text{m}^2$ Nb/Al-STJ response to VIS light at single photon level

Assuming a Poisson distribution convoluted with Gaussian which has same sigma as pedestal noise:

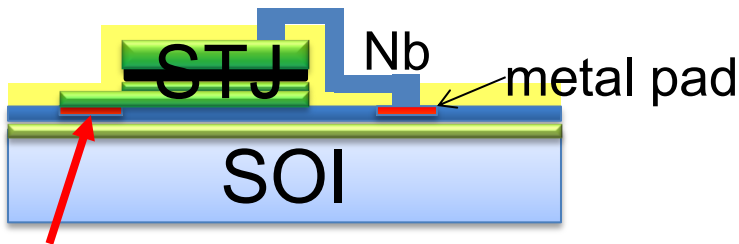
$$f(x; \mu) = N_{\text{tot}} \sum_n \frac{\mu^n e^{-\mu}}{n!} \cdot \frac{1}{\sqrt{2\pi}\sigma_p} \exp \left[ -\frac{\left\{ x - n \left( \frac{M}{\mu} \right) \right\}^2}{2\sigma_p^2} \right]$$



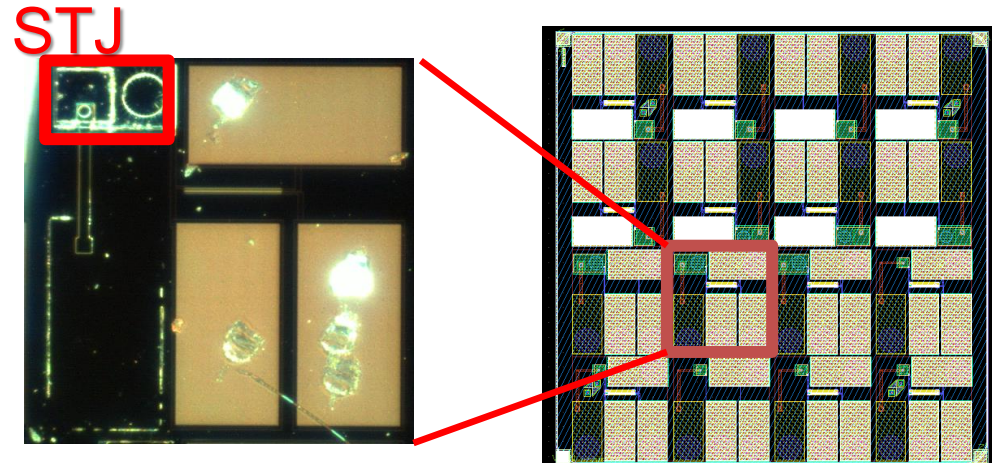
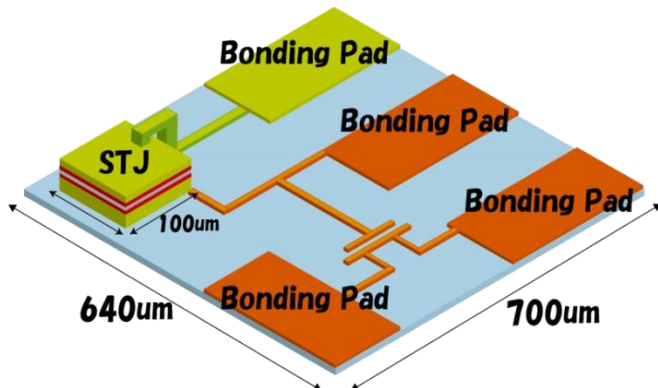
Currently, readout noise is dominated, but we are detecting VIS light at single photon level

# Development of SOI-STJ

- SOI: Silicon-on-insulator
  - CMOS in SOI is reported to work at 4.2K by T. Wada (JAXA), et al.  
Phys. 167, 602 (2012)
- A development of SOI-STJ for our application with Y. Arai (KEK)
  - STJ layer sputtered **directly on** SOI pre-amplifier
- Started test with Nb/Al-STJ on SOI with p-MOS and n-MOS FET



STJ lower layer has electrical contact with SOI circuit



**K. Kasahara's talk for detail**  
Photon detectors session3 at 15:00 on 6<sup>th</sup>

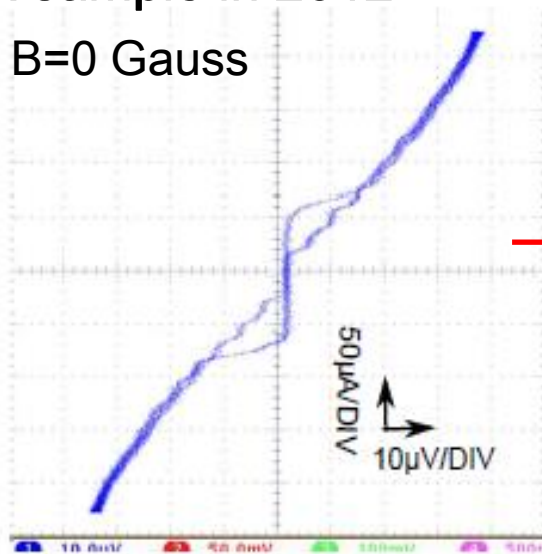


# Hf-STJ development

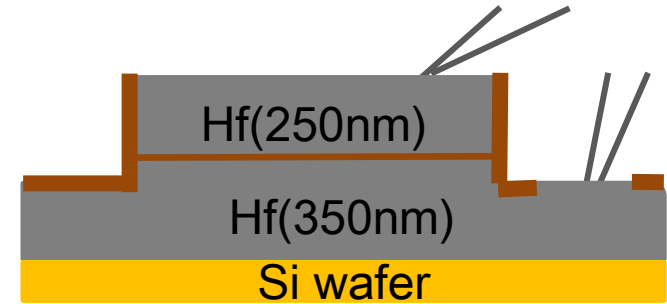
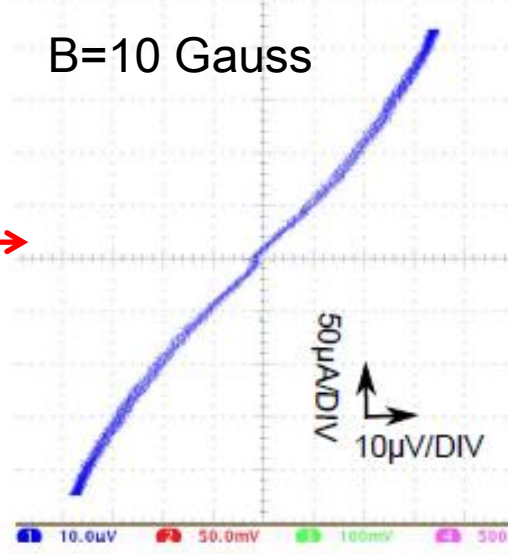
- We succeeded in observation of Josephson current by Hf-HfOx-Hf barrier layer **in 2010** (S.H.Kim et. al, TIPP2011)

## A sample in 2012

B=0 Gauss



B=10 Gauss



HfOx : 20Torr, 1hour  
anodic oxidation :  
45nm

$200 \times 200 \mu\text{m}^2$

$T = 80 \sim 177 \text{mK}$

$I_c = 60 \mu\text{A}$

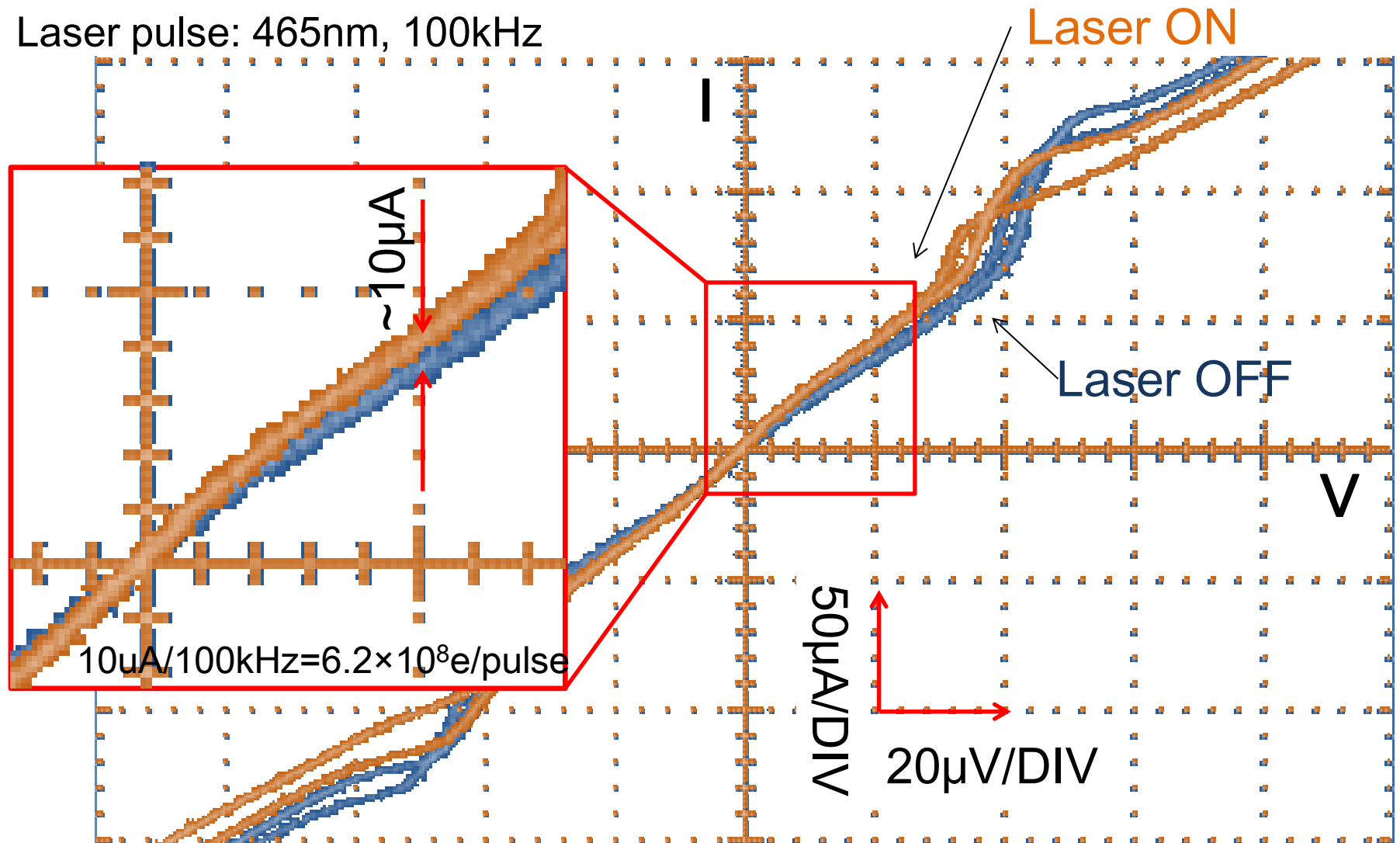
$I_{\text{leak}} = 50 \mu\text{A} @ V_{\text{bias}} = 10 \mu\text{V}$

$R_d = 0.2 \Omega$

However, to use this as a detector, much improvement in leak current is required. ( $I_{\text{leak}}$  is required to be at pA level or less)

# Hf-STJ Response to DC-like VIS light

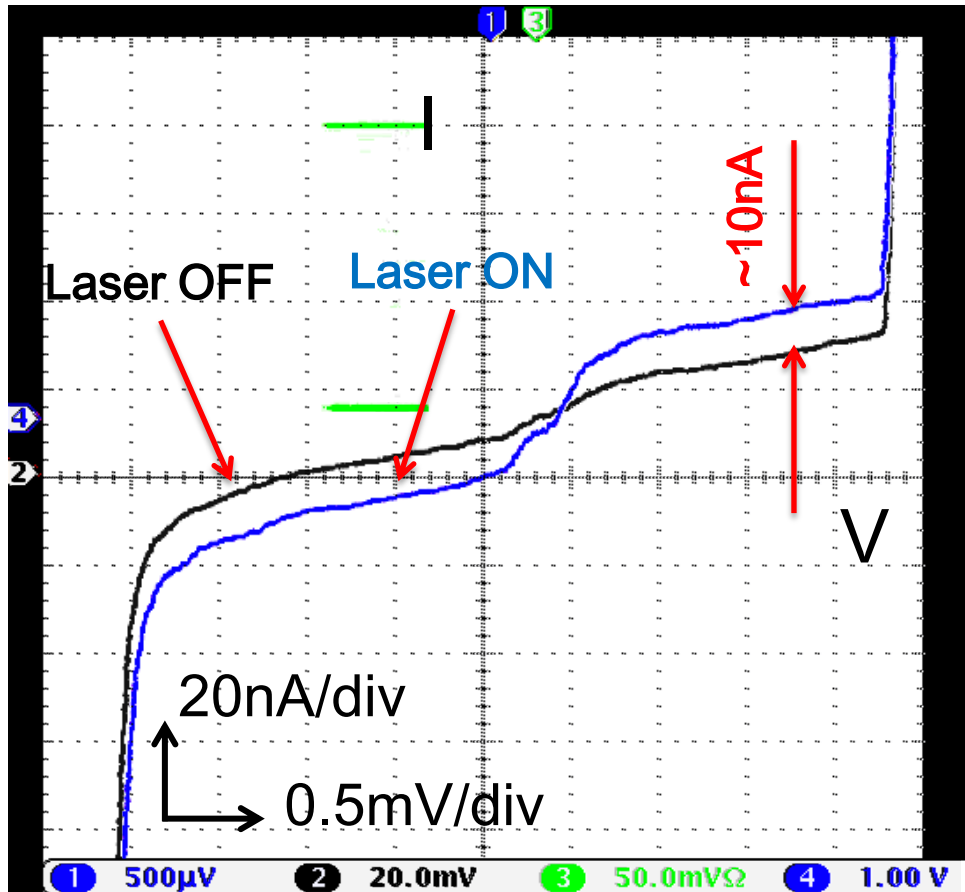
Laser pulse: 465nm, 100kHz



We observed Hf-STJ response to visible light

# 4 $\mu\text{m}^2$ Nb/Al-STJ response to DC-like VIS light

VIS laser pulse (465nm) 20MHz illuminated on 4 $\mu\text{m}^2$  Nb/Al-STJ



$$10\text{nA}/20\text{MHz}=0.5\times 10^{-15}\text{ C}=3.1\times 10^3\text{ e}$$

0.45 photons/laser pulse is given in previous slide

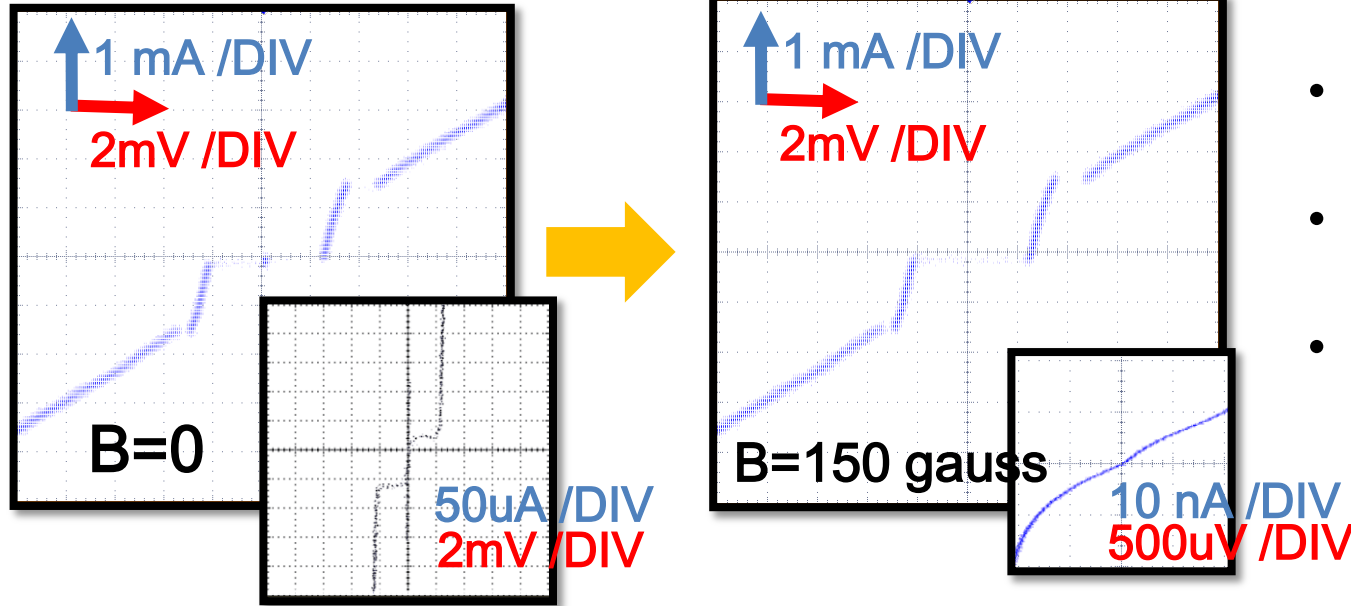


Gain from Back-tunneling effect

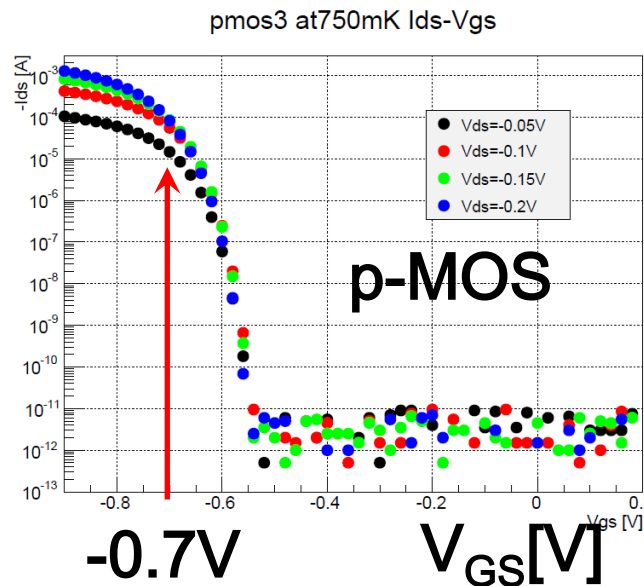
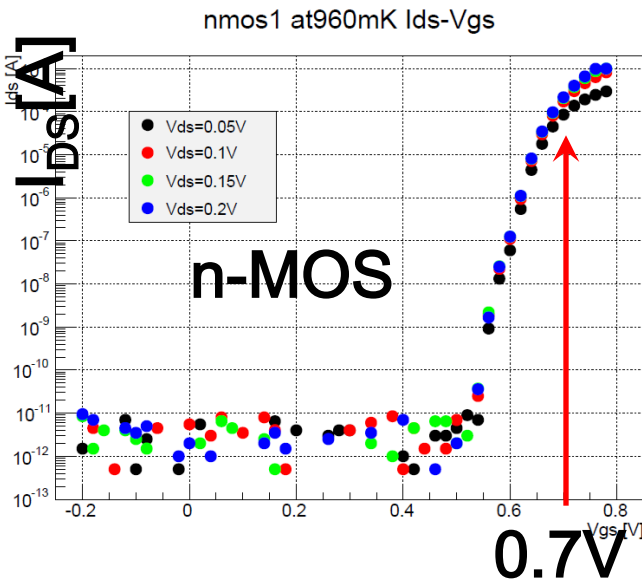
$$G_{\text{al}}=6.7$$

# Development of SOI-STJ

by K. Kasahara



- We formed Nb/Al-STJ on SOI
- Josephson current observed
- Leak current is 6nA @  $V_{\text{bias}}=0.5\text{mV}$ ,  $T=700\text{mK}$



- n-MOS and p-MOS in SOI on which STJ is formed
- Both n-MOS and p-MOS works at  $T=750\sim 960\text{mK}$

# STJ on SOI response to laser pulse

Illuminate 20 laser pulses (465nm, 50MHz) on 50x50um<sup>2</sup> STJ which is formed on SOI

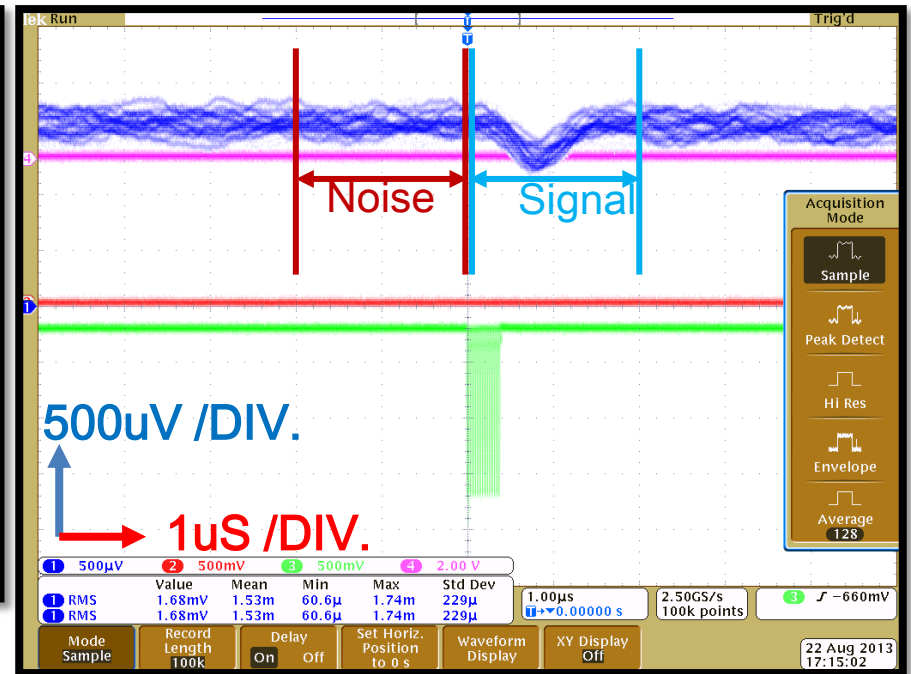
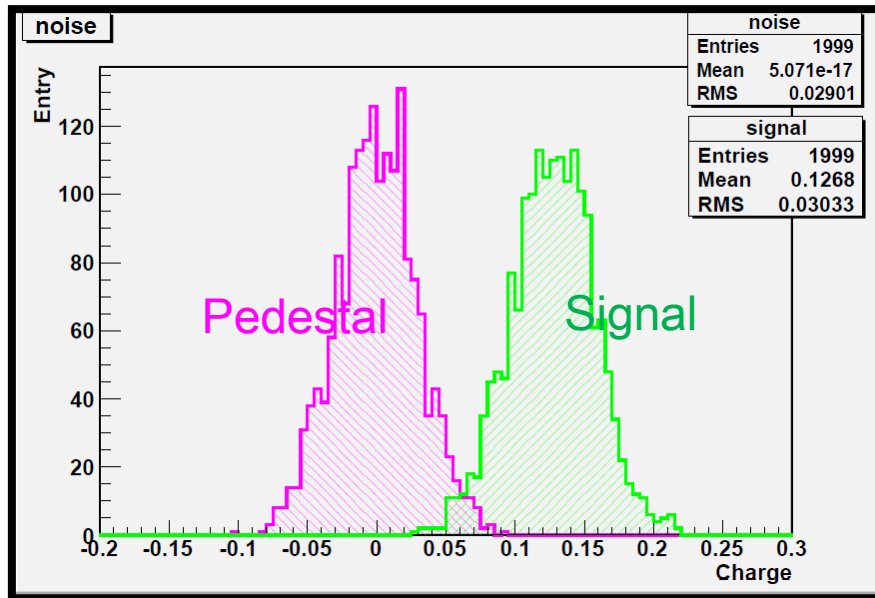
Estimated number of photons from output signal pulse height distribution assuming photon statistics

$$N_Y = \frac{M^2}{\sigma^2 - \sigma_p^2} \sim 206 \pm 112$$

M : Mean

$\sigma$  : signal RMS

$\sigma_p$  : pedestal RMS



Confirmed STJ formed on SOI responds to VIS laser pulses